

Gravitational Waves: Theoretical Insight to Measurement

Rainer Weiss, MIT

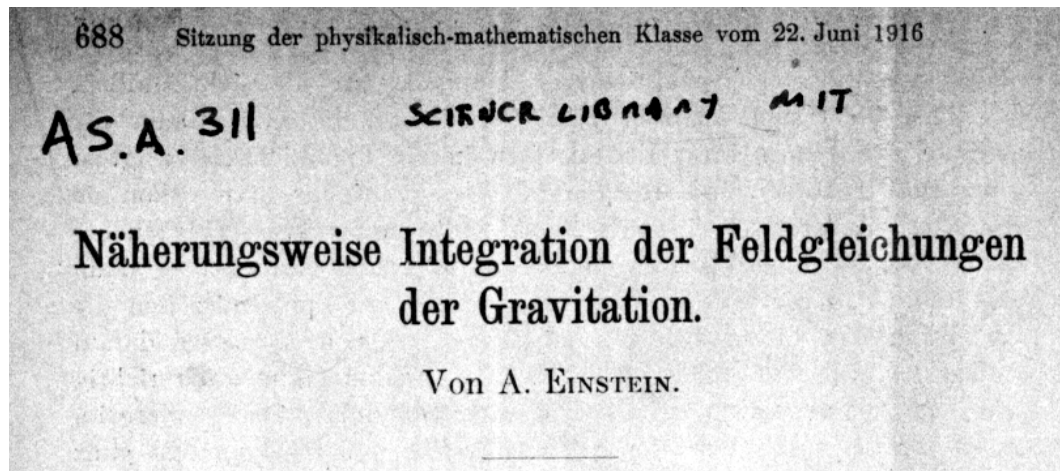
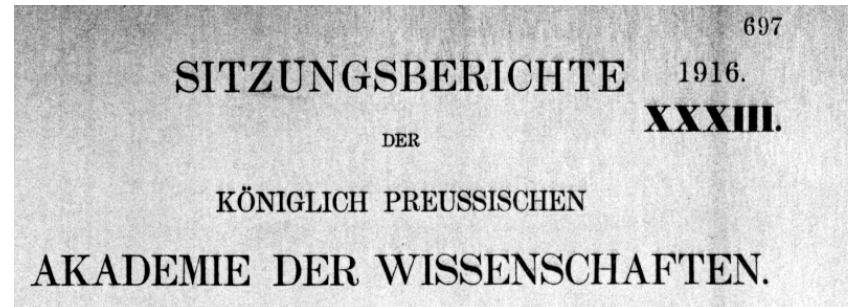
A Century of General Relativity

Harnack House

Berlin

November 30, 2015

Einstein 1916



Einstein 1916

$$\gamma'_{\mu\nu} = \alpha_{\mu\nu} f(x_1 + ix_4) = \alpha_{\mu\nu} f(x-t). \quad (15)$$

Dabei sind die $\alpha_{\mu\nu}$ Konstante; f ist eine Funktion des Arguments $x-t$. Ist der betrachtete Raum frei von Materie, d. h. verschwinden die $T_{\mu\nu}$, so sind die Gleichungen (6) durch diesen Ansatz erfüllt. Die Gleichungen (4) liefern zwischen den $\alpha_{\mu\nu}$ die Beziehungen

$$\left. \begin{aligned} \alpha_{11} + i\alpha_{14} &= 0 \\ \alpha_{12} + i\alpha_{24} &= 0 \\ \alpha_{13} + i\alpha_{34} &= 0 \\ \alpha_{14} + i\alpha_{44} &= 0 \end{aligned} \right\}. \quad (16)$$

Von den 10 Konstanten $\alpha_{\mu\nu}$ sind daher nur 6 frei wählbar. Wir können die allgemeinste Welle der betrachteten Art daher aus Wellen von folgenden 6 Typen superponieren

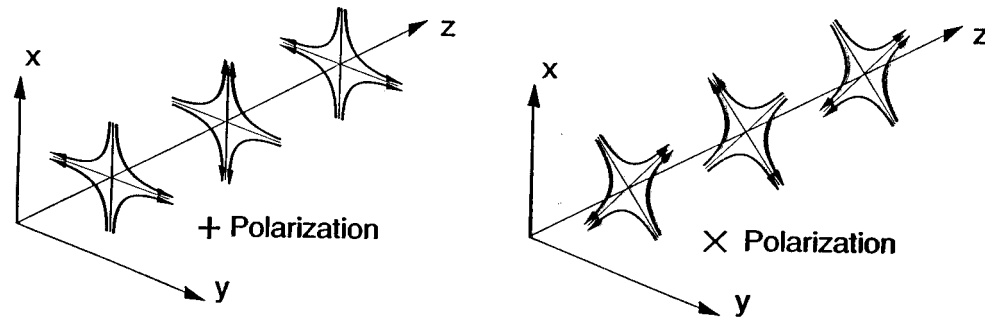
$$\left. \begin{array}{lll} \text{a) } \alpha_{11} + i\alpha_{14} = 0 & \text{b) } \alpha_{12} + i\alpha_{24} = 0 & \text{d) } \alpha_{22} \neq 0 \\ \alpha_{14} + i\alpha_{44} = 0 & \text{c) } \alpha_{13} + i\alpha_{34} = 0 & \text{e) } \alpha_{23} \neq 0 \\ & & \text{f) } \alpha_{33} \neq 0 \end{array} \right\}. \quad (17)$$

$$\begin{aligned} \text{d) } \frac{1}{i} t_{22} &= \frac{f'^2}{4x} \alpha_{22}^2 = \frac{1}{4x} \left(\frac{\partial \gamma'_{22}}{\partial t} \right)^2 \\ \text{e) } \frac{1}{i} t_{23} &= \frac{f'^2}{4x} \alpha_{23}^2 = \frac{1}{4x} \left(\frac{\partial \gamma'_{23}}{\partial t} \right)^2 \\ \text{f) } \frac{1}{i} t_{33} &= \frac{f'^2}{4x} \alpha_{33}^2 = \frac{1}{4x} \left(\frac{\partial \gamma'_{33}}{\partial t} \right)^2 \end{aligned}$$

Es ergibt sich also, daß nur die Wellen des letzten Typs Energie transportieren, und zwar ist der Energietransport einer beliebigen ebenen Welle gegeben durch

$$I_x = \frac{1}{i} t_{41} = \frac{1}{4x} \left[\left(\frac{\partial \gamma'_{22}}{\partial t} \right)^2 + 2 \left(\frac{\partial \gamma'_{23}}{\partial t} \right)^2 + \left(\frac{\partial \gamma'_{33}}{\partial t} \right)^2 \right]. \quad (18)$$

Einstein 1916



Die in (23), (23a) und (23b) auftretenden Integrale, welche nichts anderes sind als zeitlich variable Trägheitsmomente, nennen wir im folgenden zur Abkürzung J_{22} , J_{33} , J_{23} . Dann ergibt sich für die Intensität \dot{I}_r der Energiestrahlung aus (18)

$$\dot{I}_r = \frac{\kappa}{64\pi^2 R^2} \left[\left(\frac{\partial^3 J_{22}}{\partial t^3} \right)^2 + 2 \left(\frac{\partial^3 J_{23}}{\partial t^3} \right)^2 + \left(\frac{\partial^3 J_{33}}{\partial t^3} \right)^2 \right]. \quad (20)$$

SPHERICALLY SYMMETRIC MOTION RADIATES GRAVITATIONAL WAVES

Einstein 2018

154 Gesamtsitzung vom 14. Februar 1918. — Mitteilung vom 31. Januar

Über Gravitationswellen.

VON A. EINSTEIN.

(Vorgelegt am 31. Januar 1918 [s. oben S. 79].)

Die wichtige Frage, wie die Ausbreitung der Gravitationsfelder erfolgt, ist schon vor anderthalb Jahren in einer Akademiearbeit von mir behandelt worden¹. Da aber meine damalige Darstellung des Gegenstandes nicht genügend durchsichtig und außerdem durch einen bedauerlichen Rechenfehler verunstaltet ist, muß ich hier nochmals auf die Angelegenheit zurückkommen.

Die $\mathfrak{J}_{\alpha\gamma}$ sind nach (7 a), (22), (24) für die Zeit $t-R$ zu nehmen, also Funktionen von $t-R$, oder bei großem R in der Nähe der x -Achse auch Funktionen von $t-x$. (25), (26) stellen also Gravitationswellen dar, deren Energiefluß längs der x -Achse gemäß (16) die Dichte

$$\frac{t_{41}}{i} = \frac{\kappa}{64 \pi^2 R^2} \left[\left(\ddot{\mathfrak{J}}_{22} - \ddot{\mathfrak{J}}_{33} \right)^2 + \dot{\mathfrak{J}}_{23}^2 \right] \quad (27)$$

besitzt.

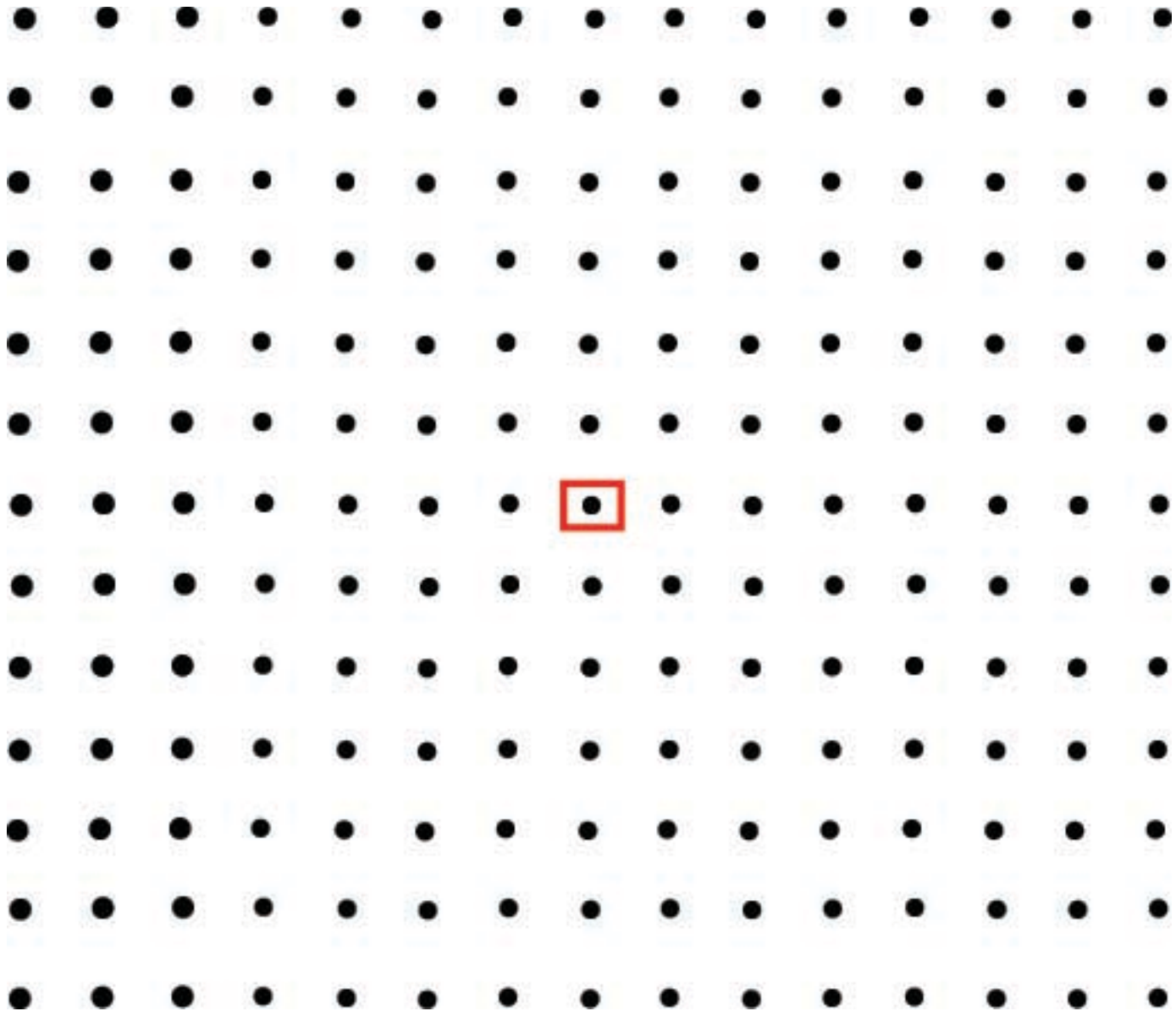
The quadrupole formula a factor of 2 too small

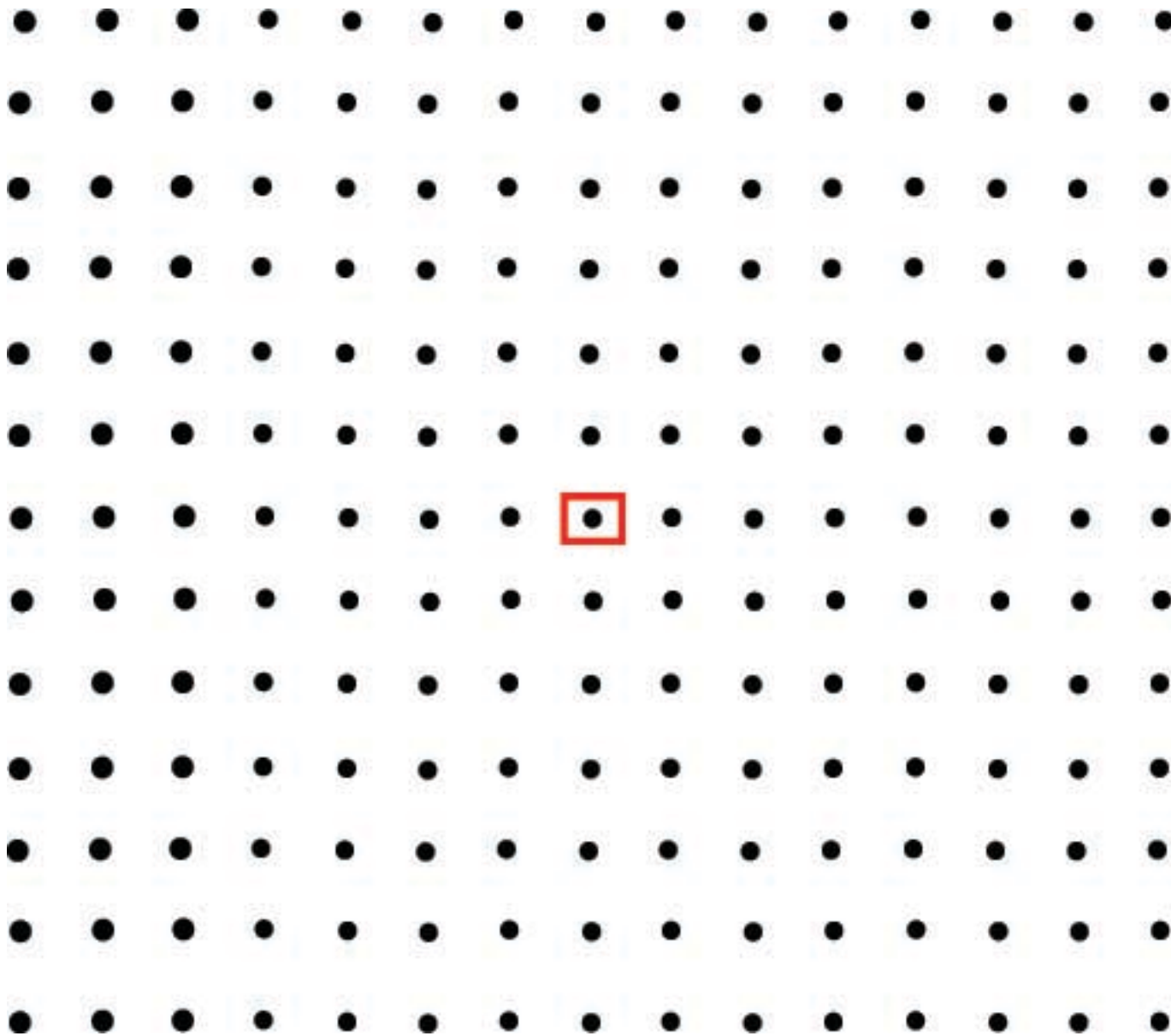
Einstein 1916

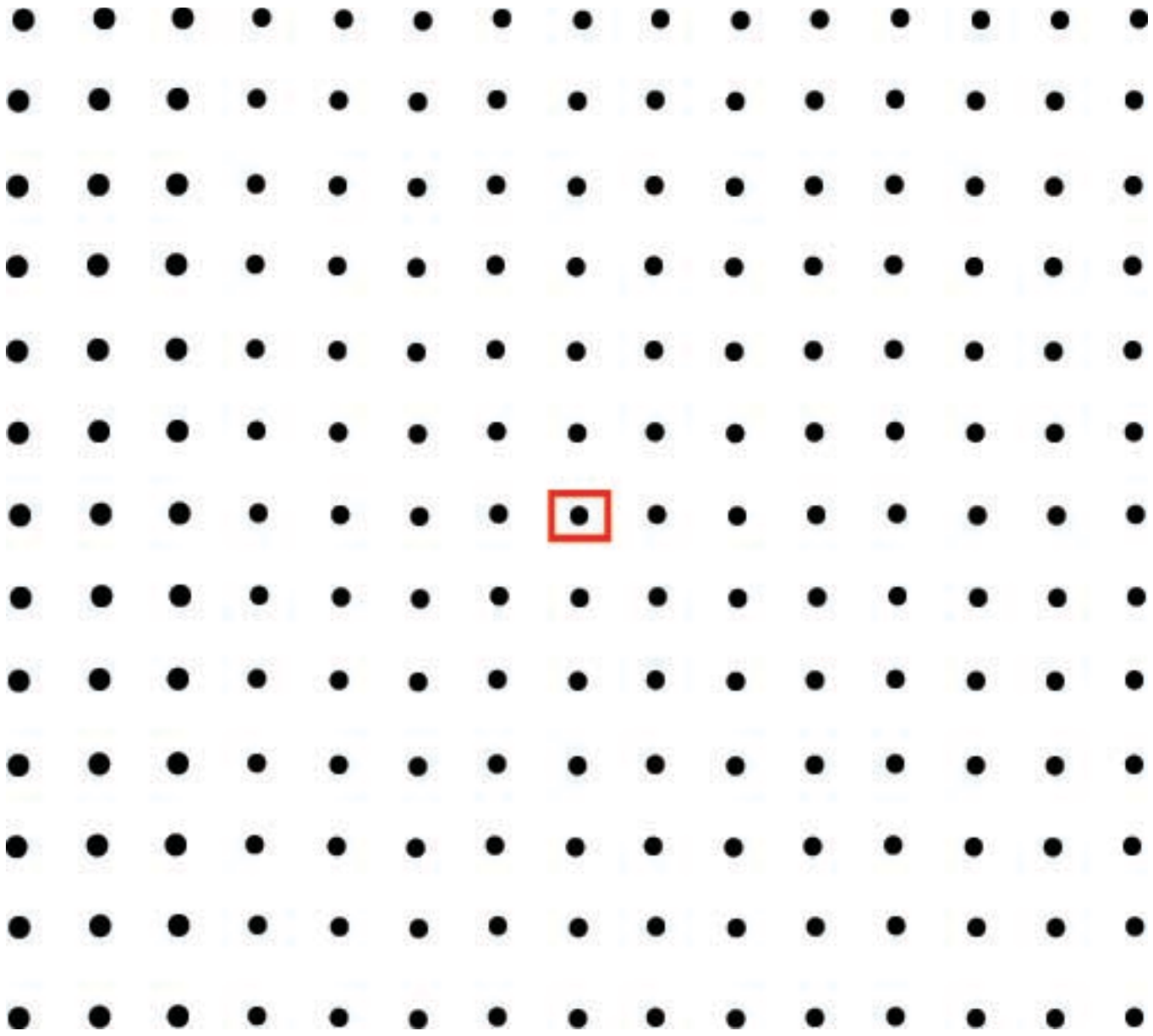
$$A = \frac{\kappa}{24\pi} \sum_{\alpha\beta} \left(\frac{\partial^3 J_{\alpha\beta}}{\partial t^3} \right)^2. \quad (21)$$

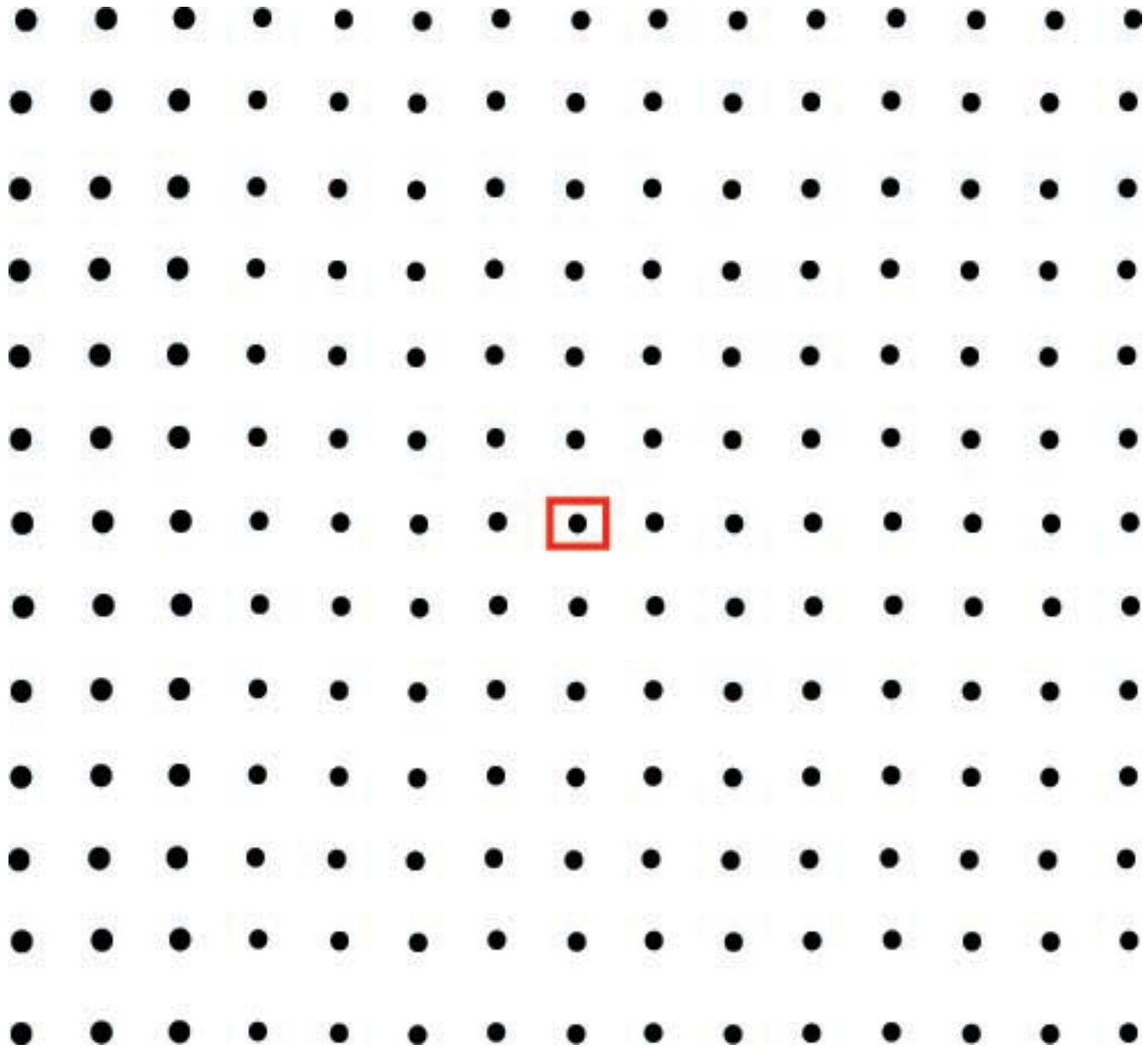
Würde man die Zeit in Sekunden, die Energie in Erg messen, so würde zu diesem Ausdruck der Zahlenfaktor $\frac{1}{c^4}$ hinzutreten. Berücksichtigt man außerdem, daß $\kappa = 1.87 \cdot 10^{-27}$, so sieht man, daß A in allen nur denkbaren Fällen einen praktisch verschwindenden Wert haben muß.

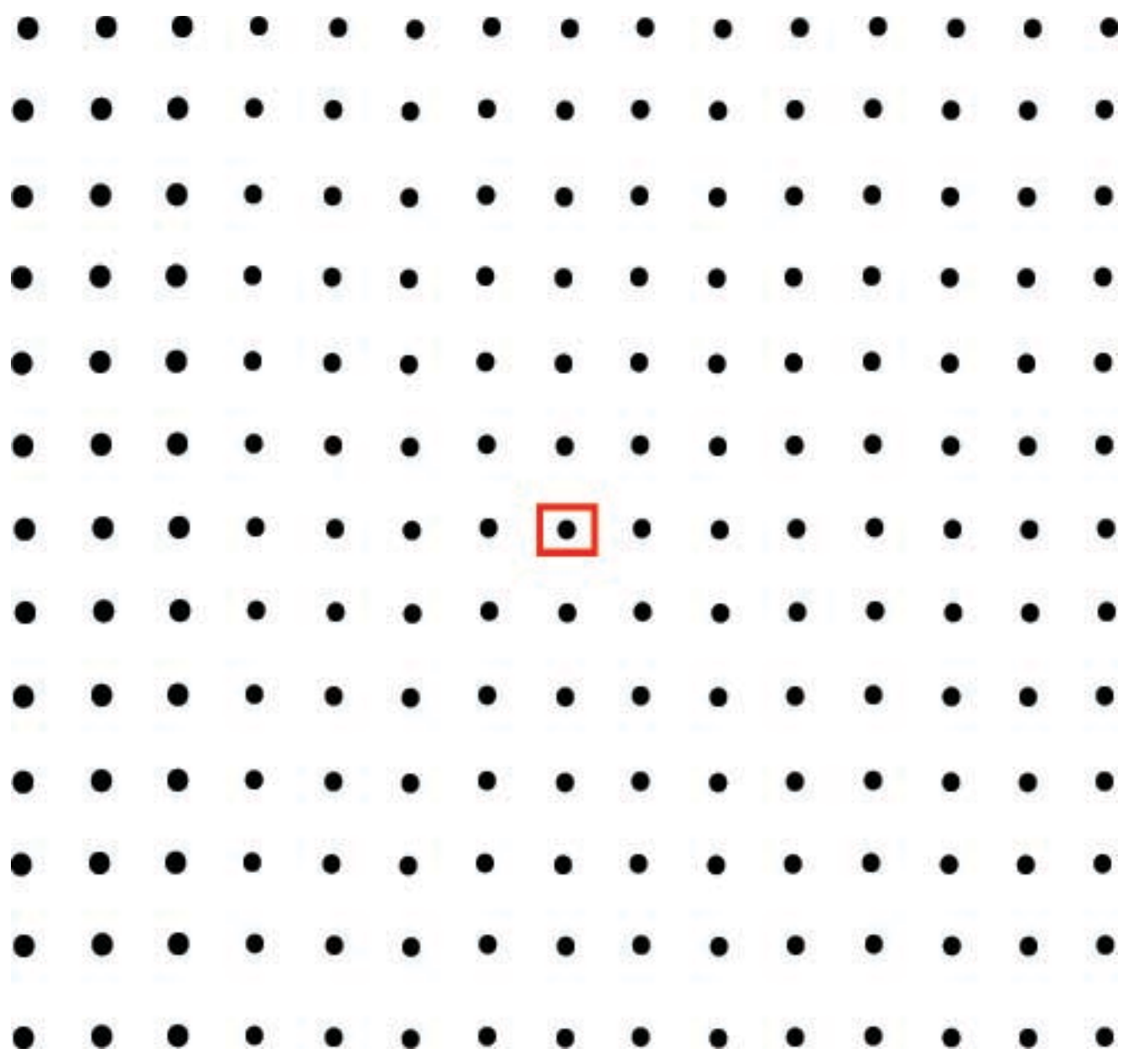
“.....in any case one can think of A will have a practically vanishing value.”

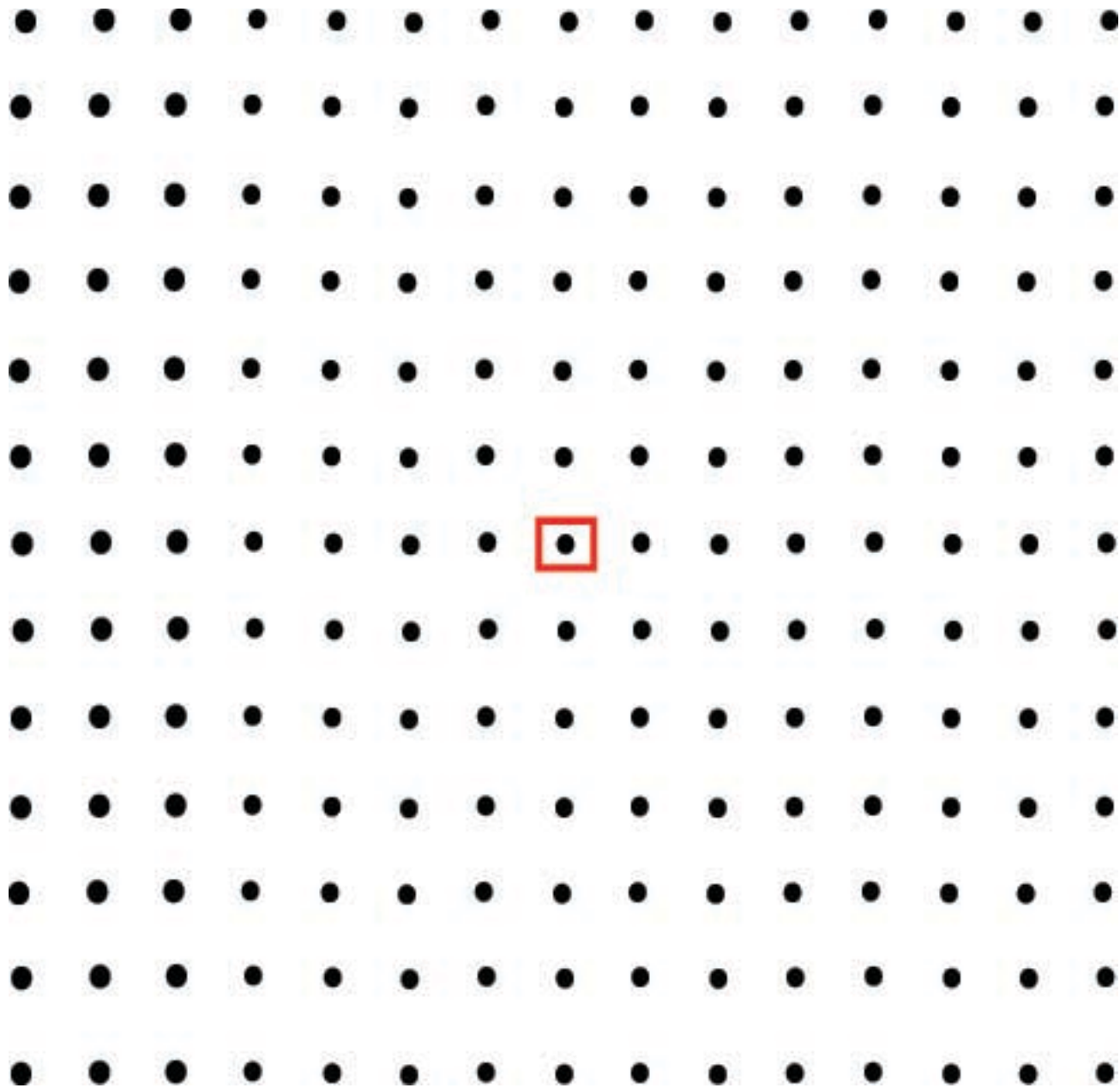


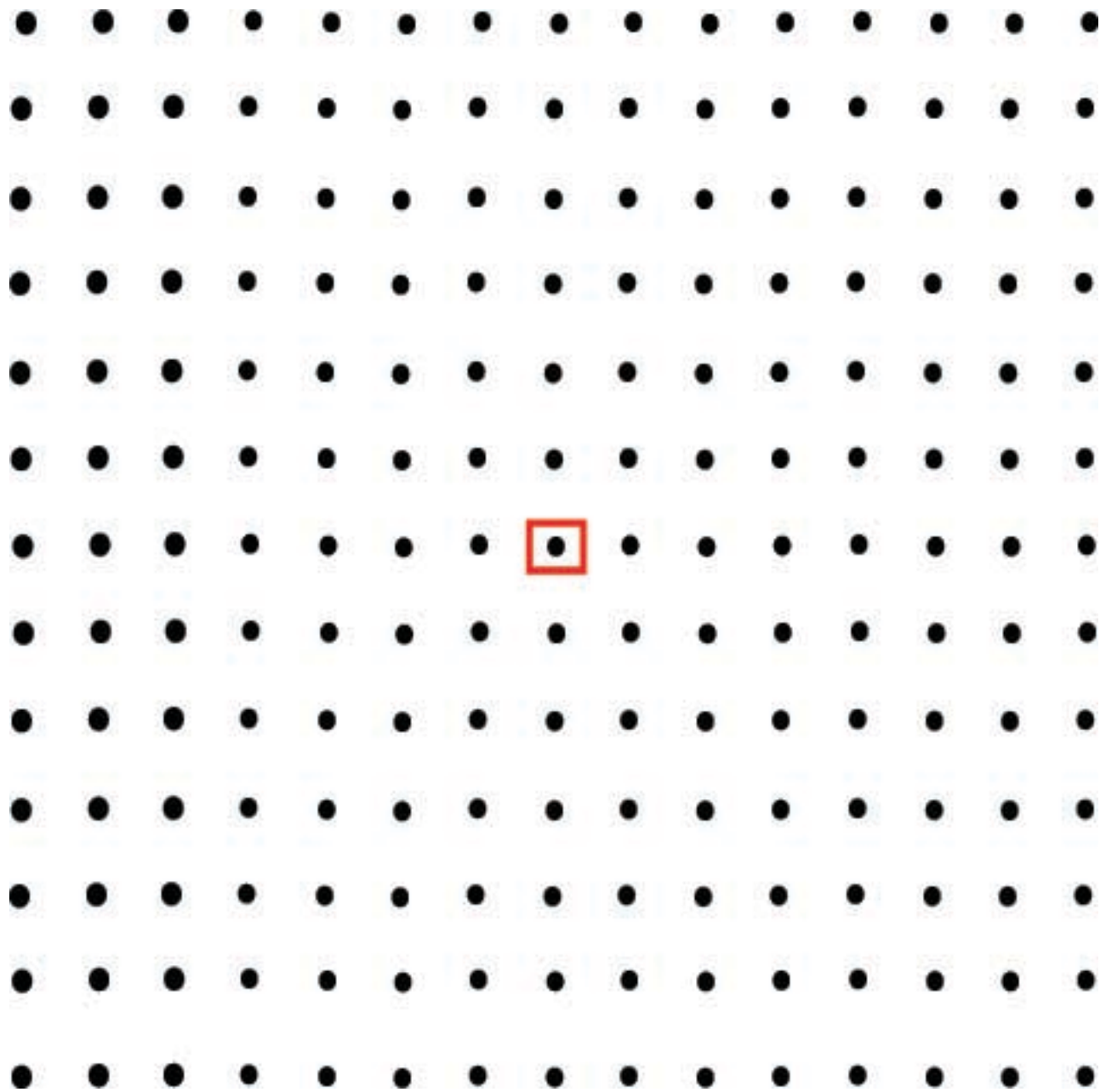


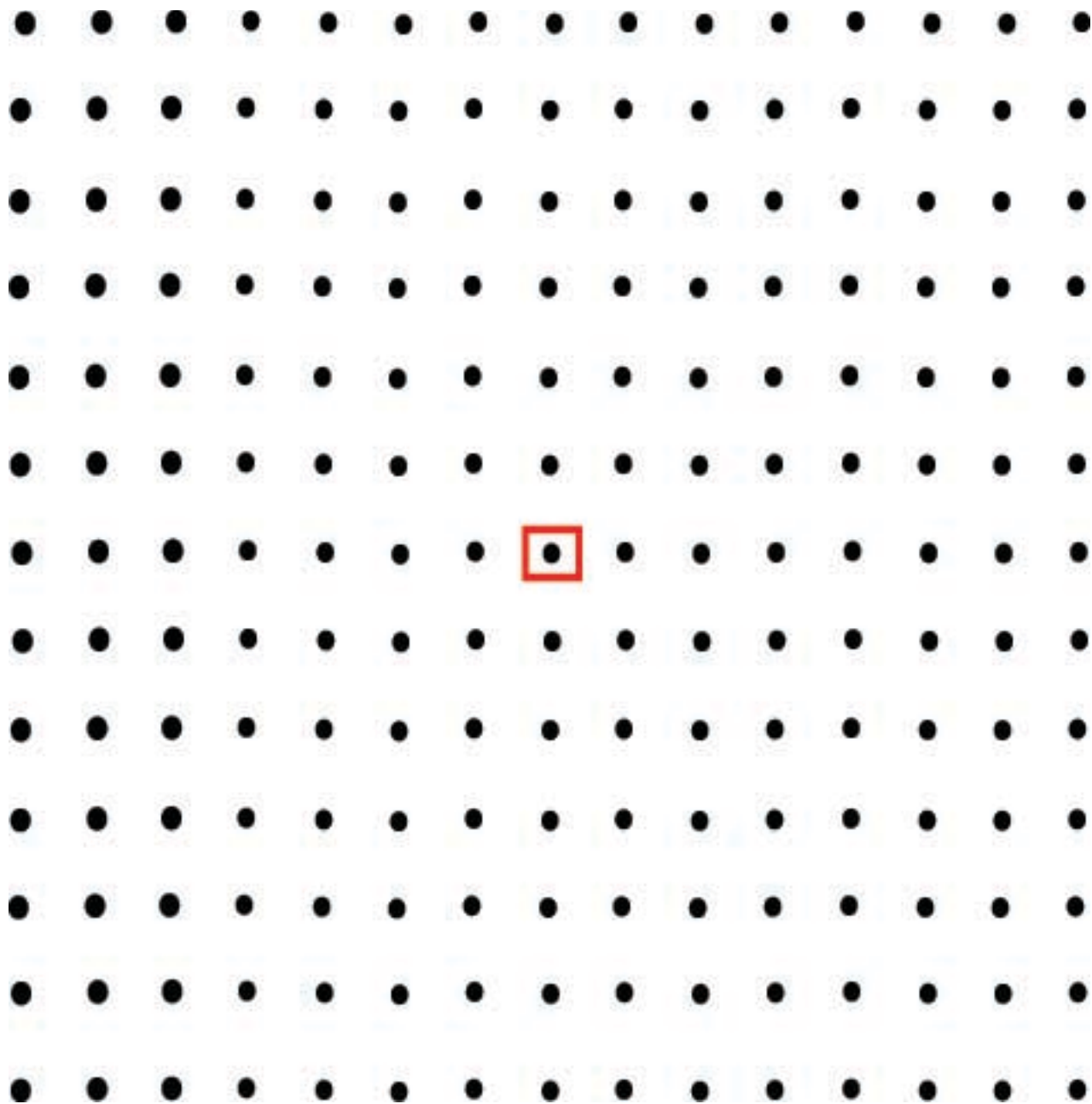


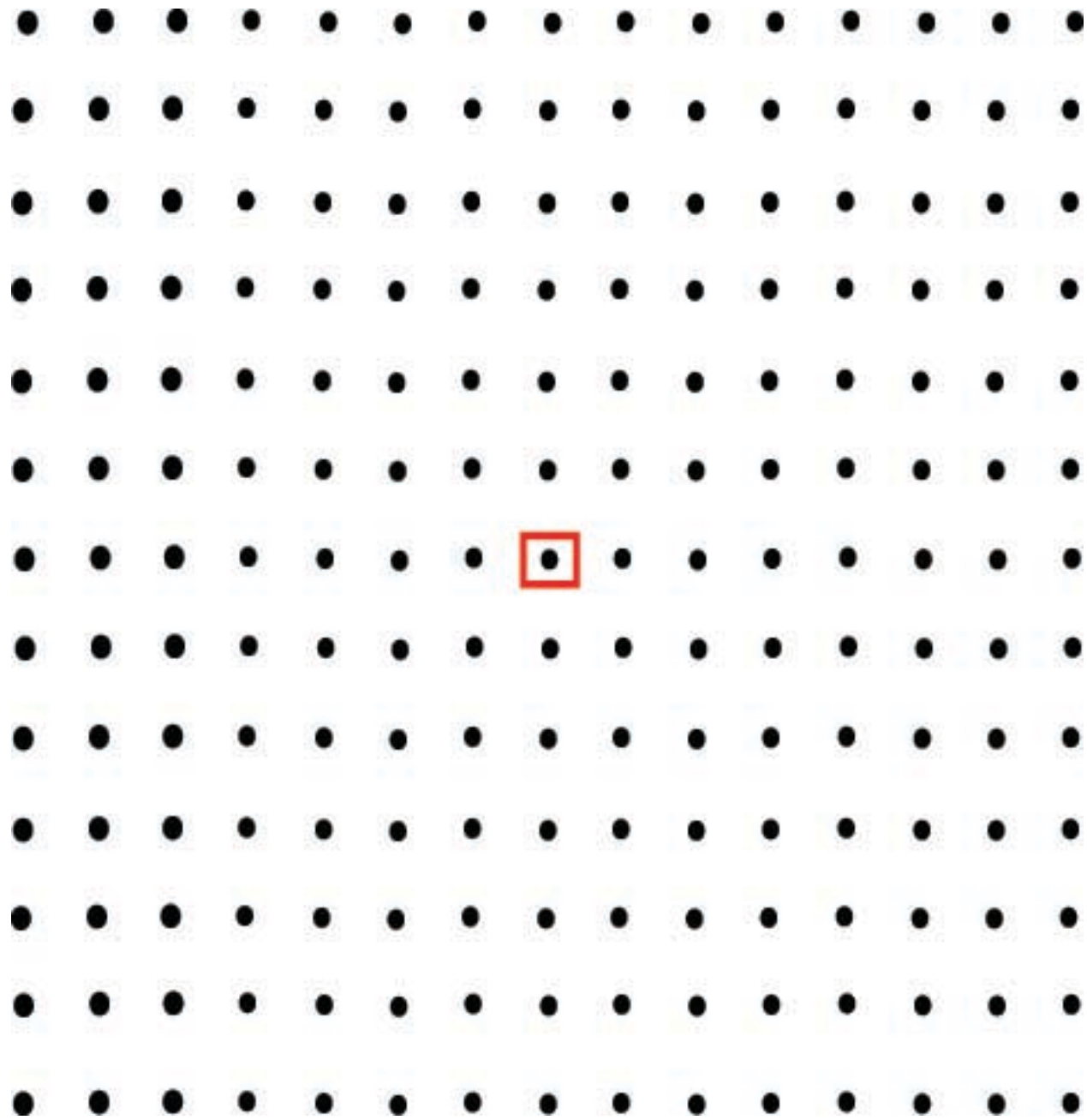


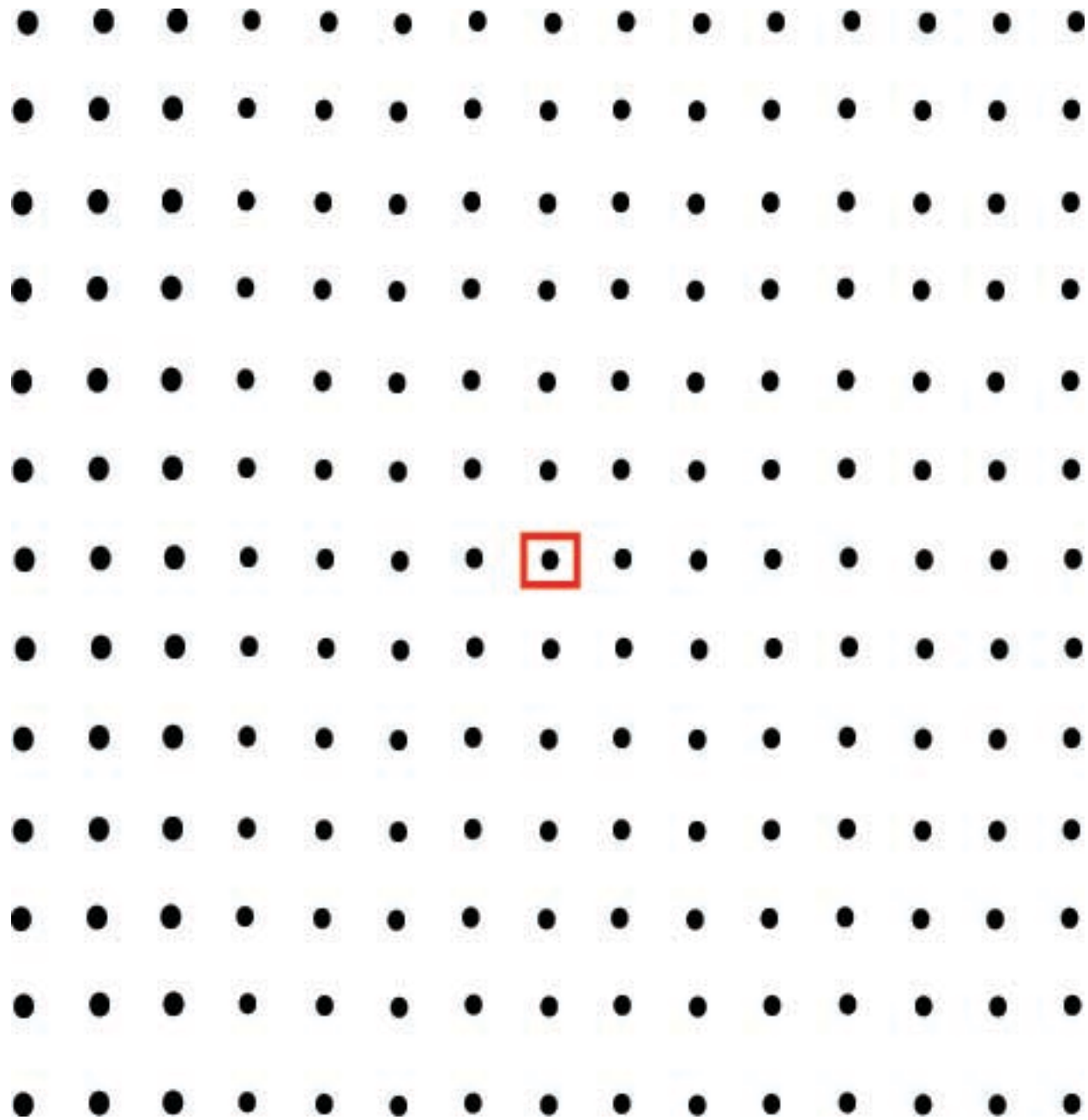


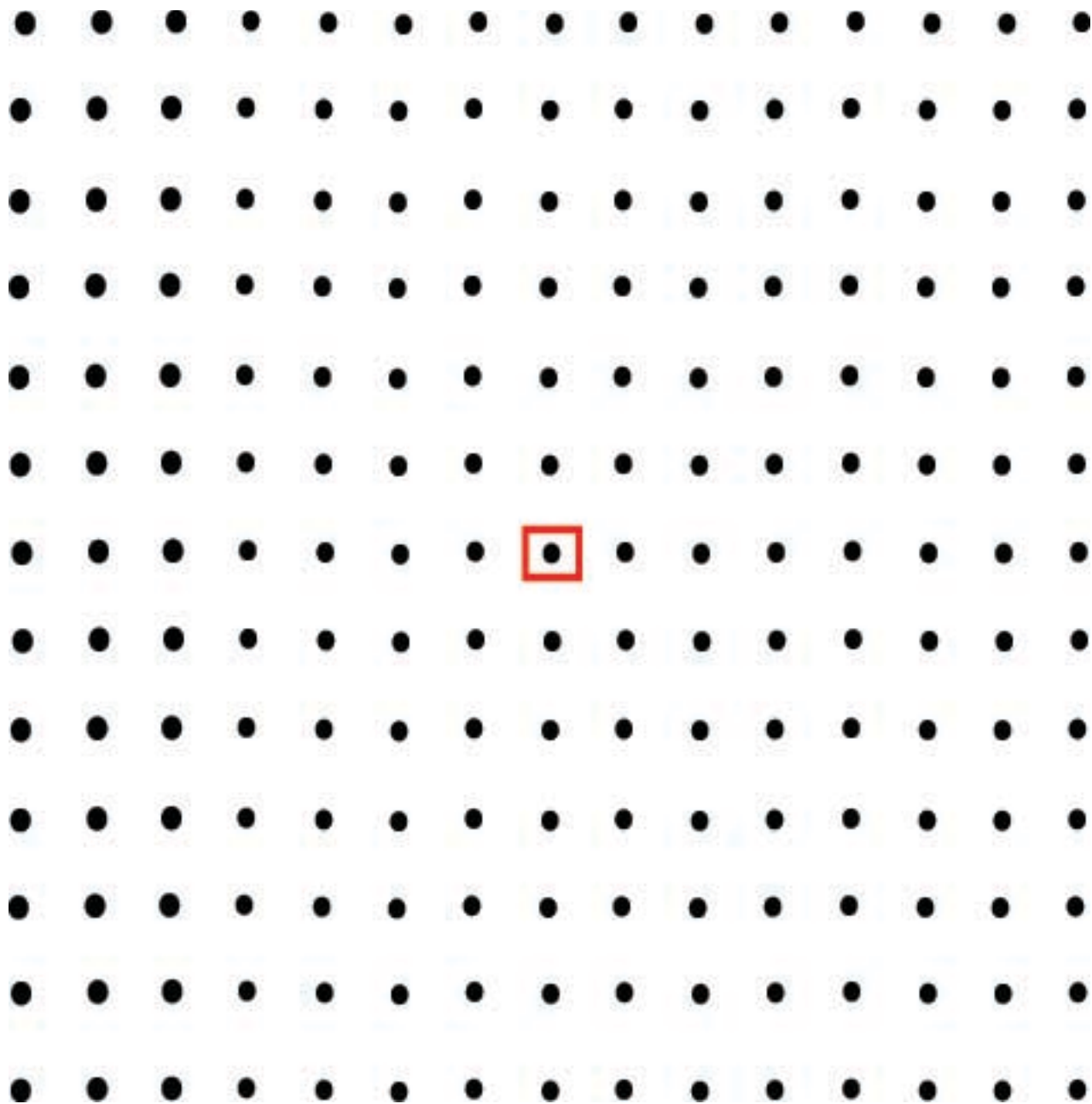


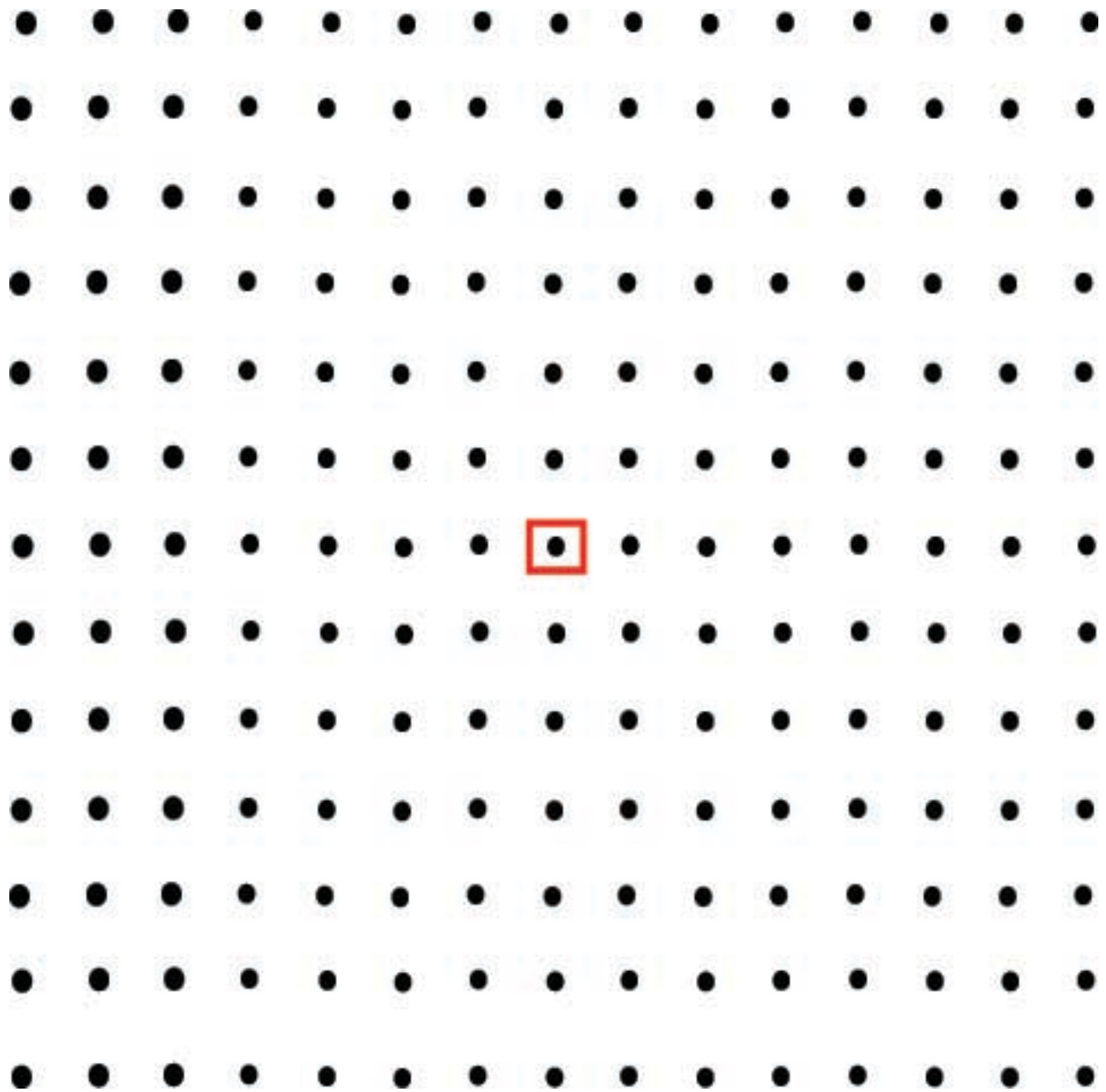


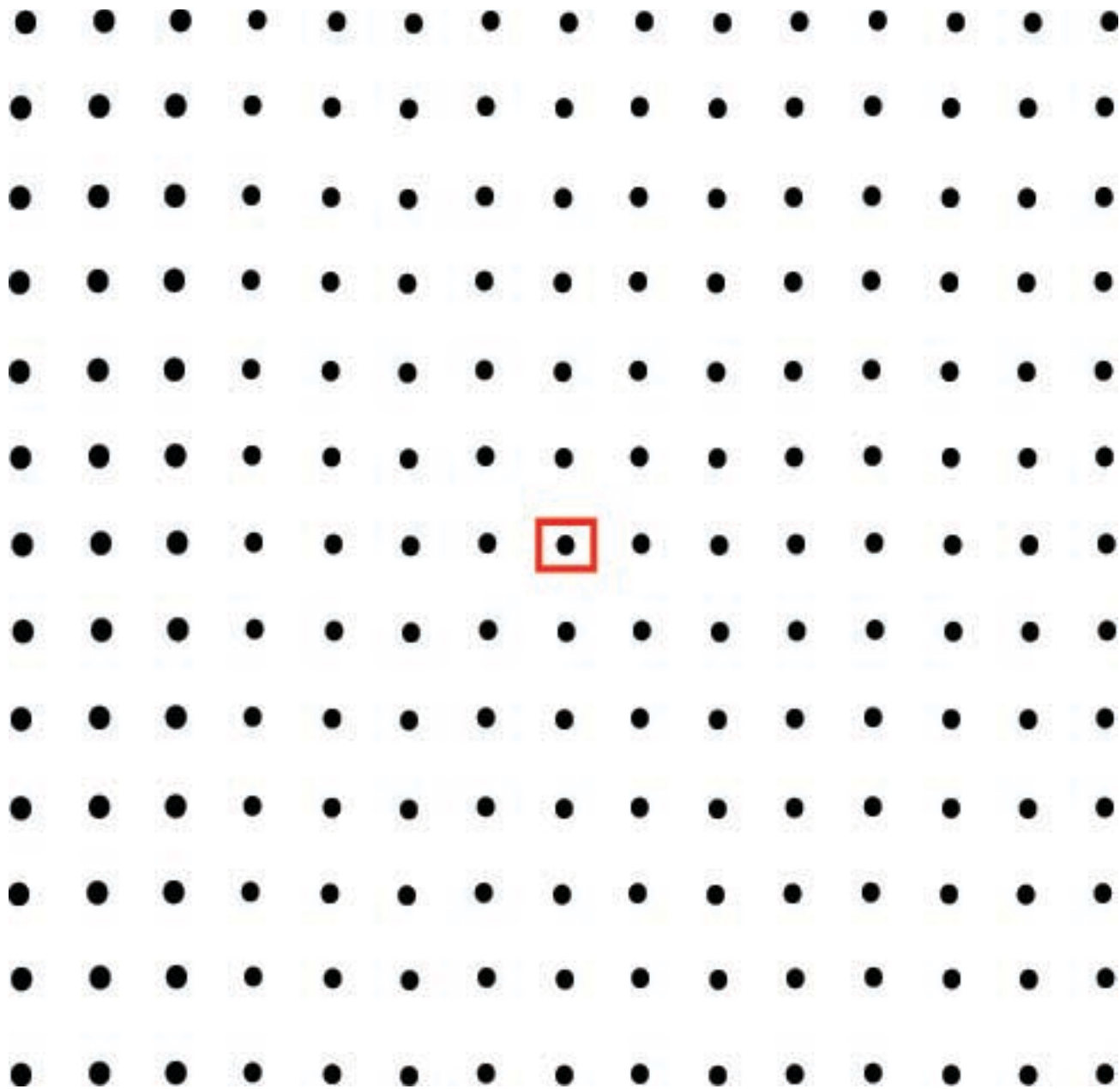


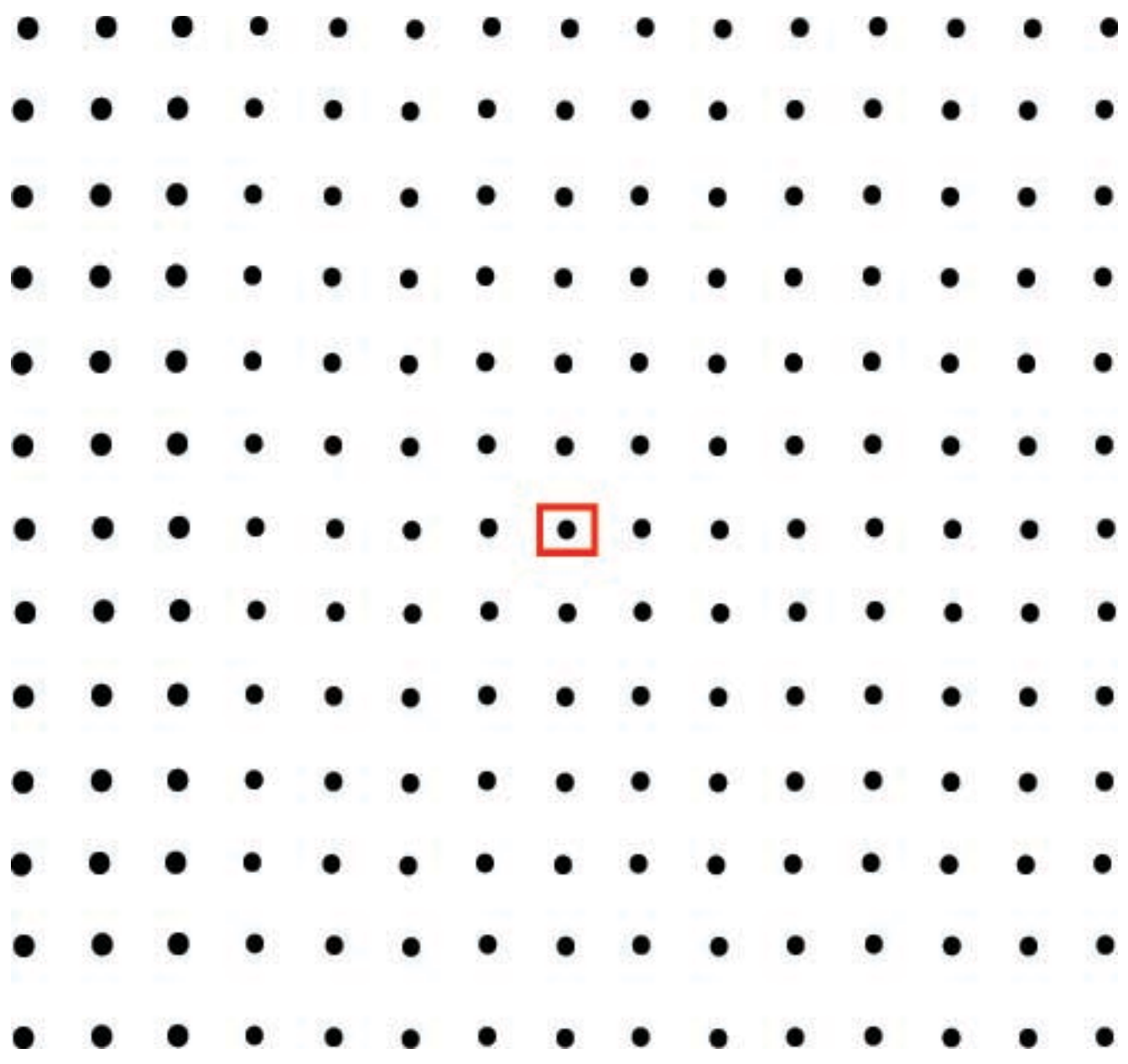


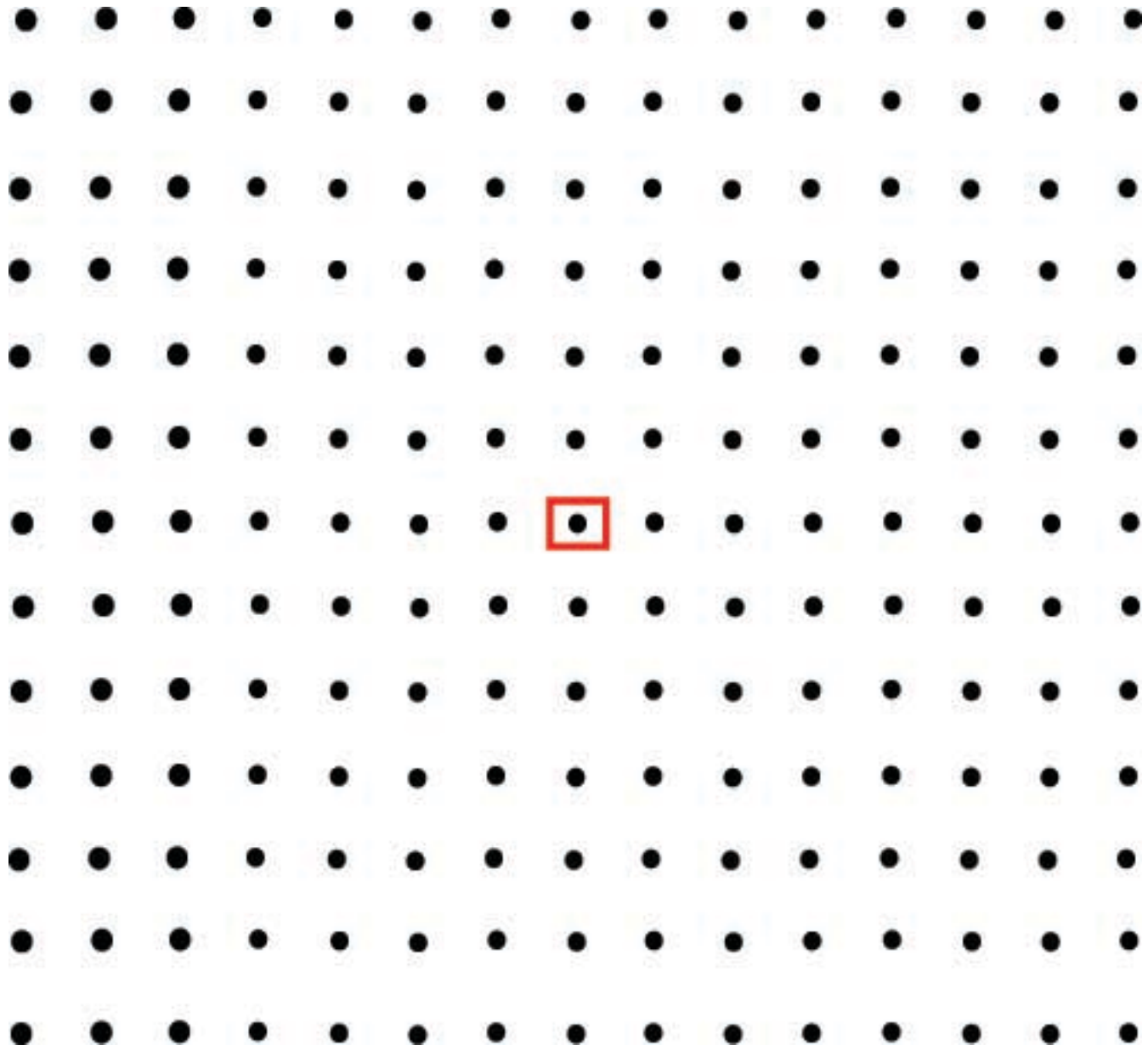


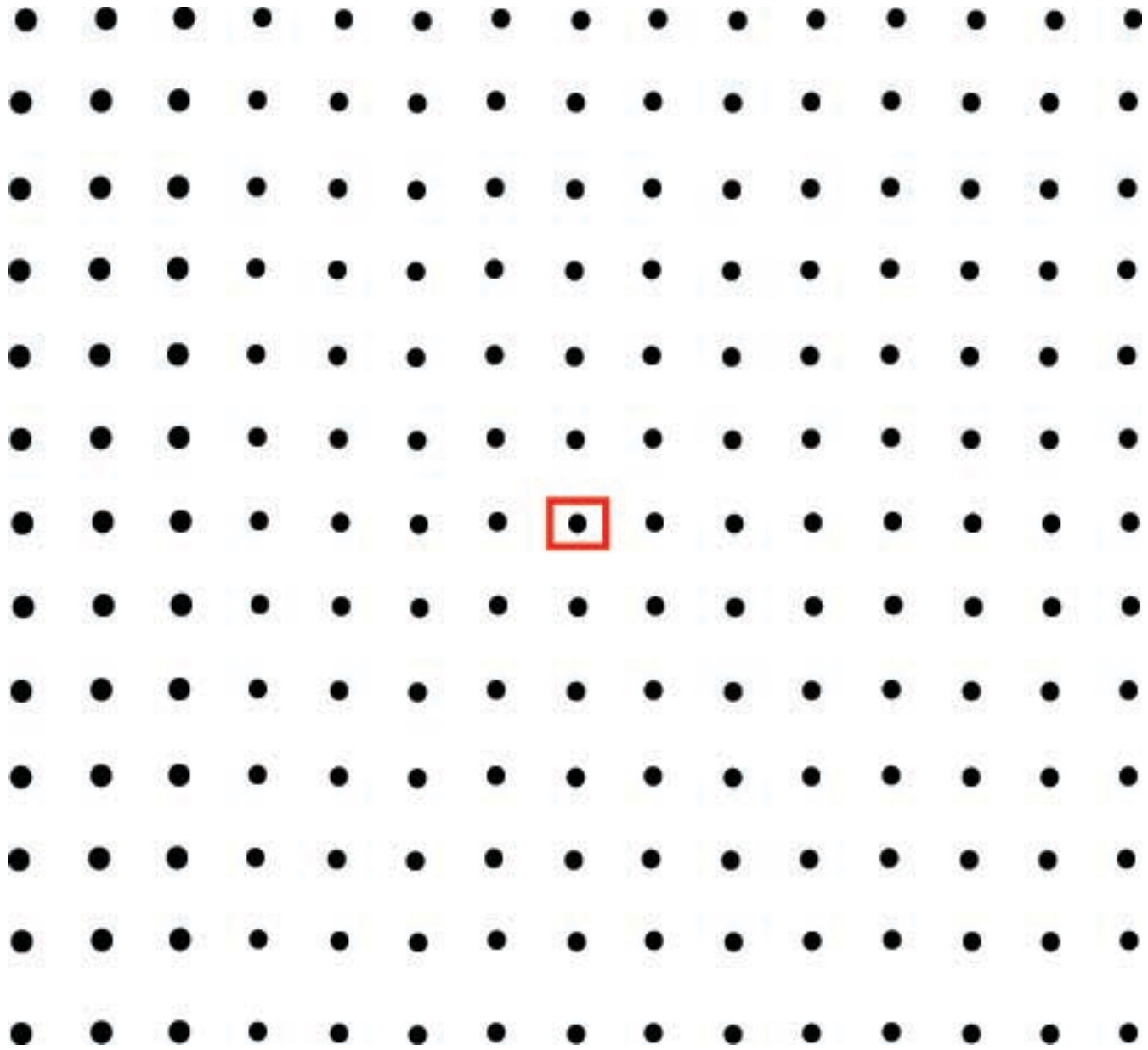


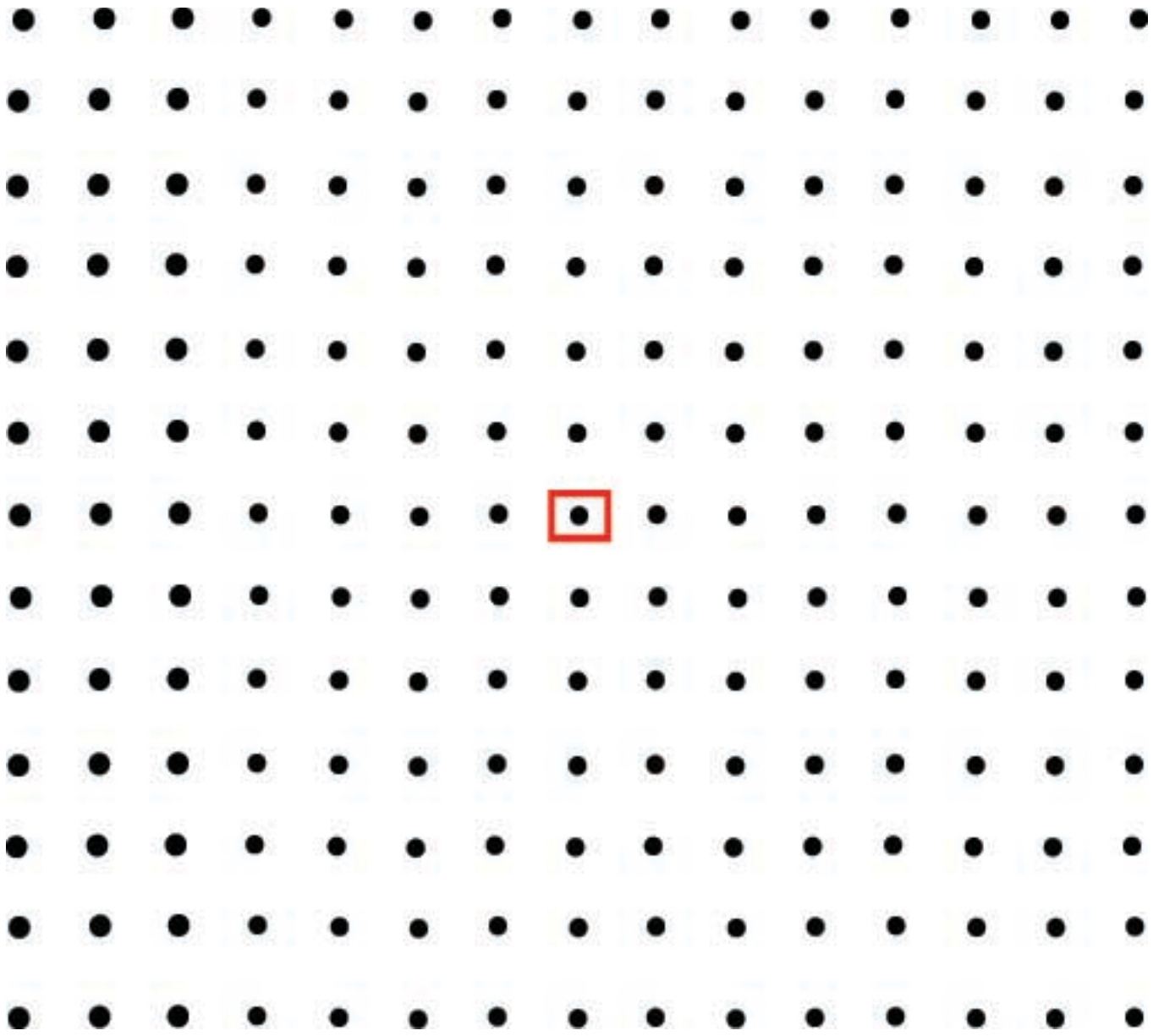


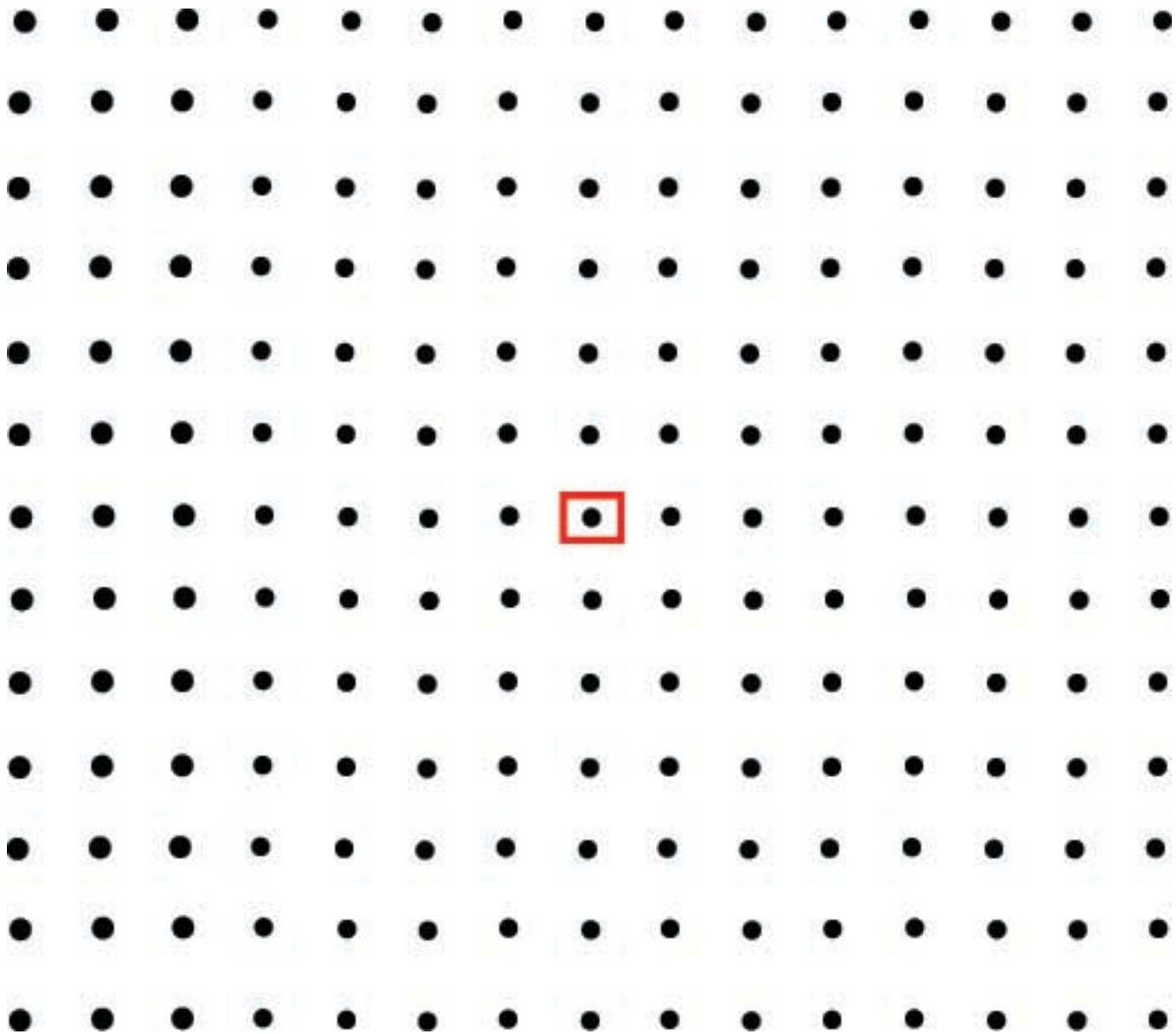


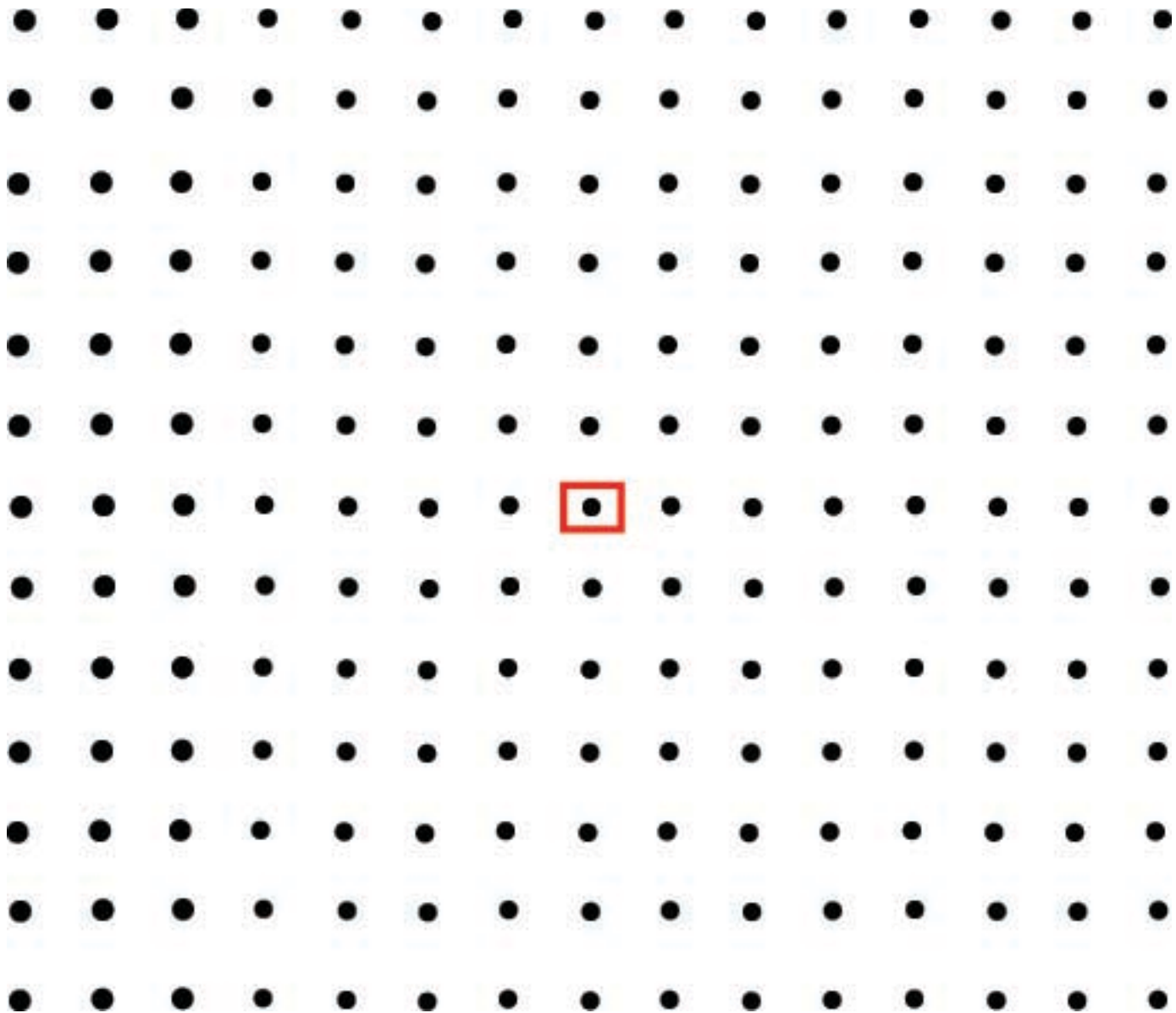


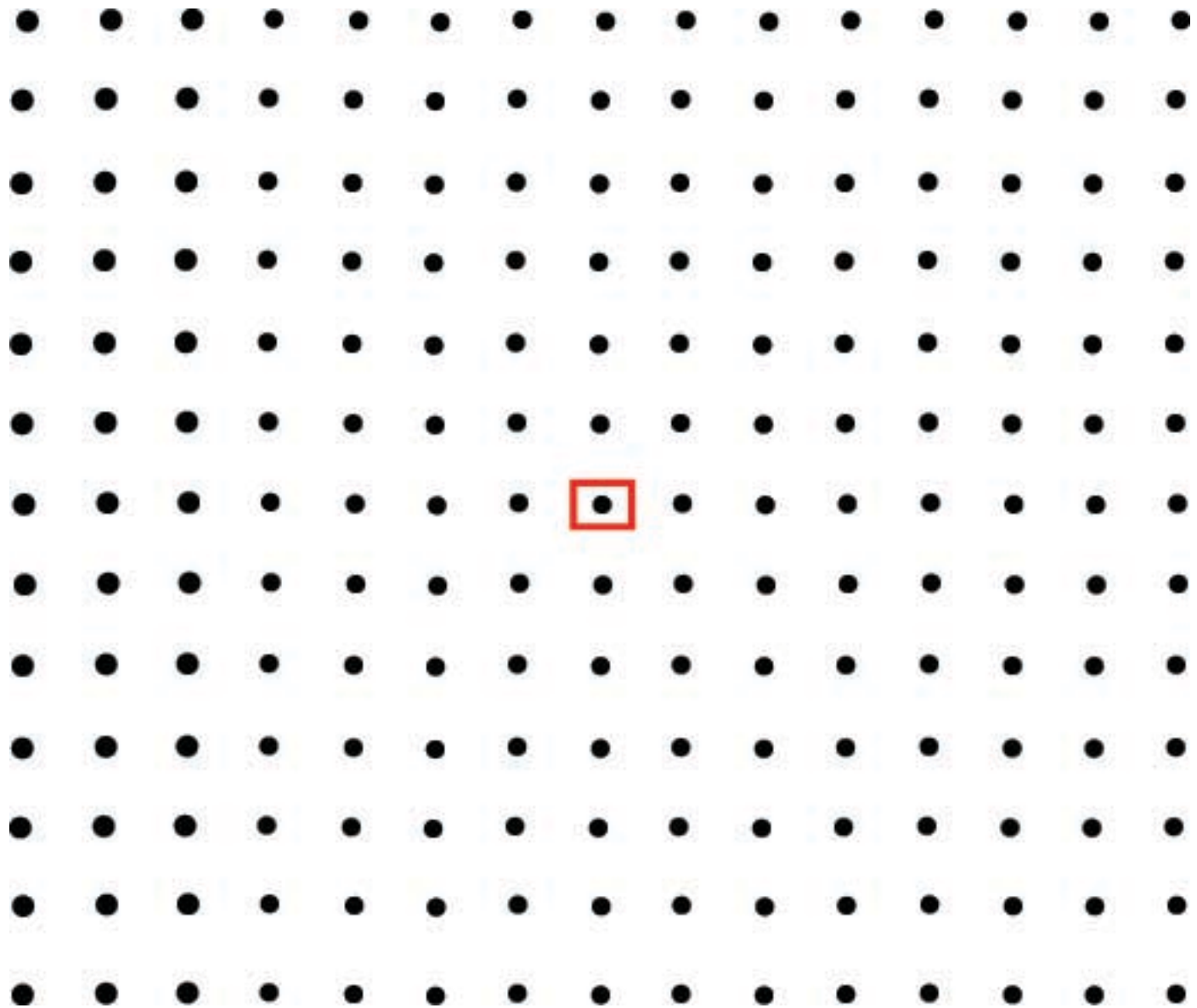


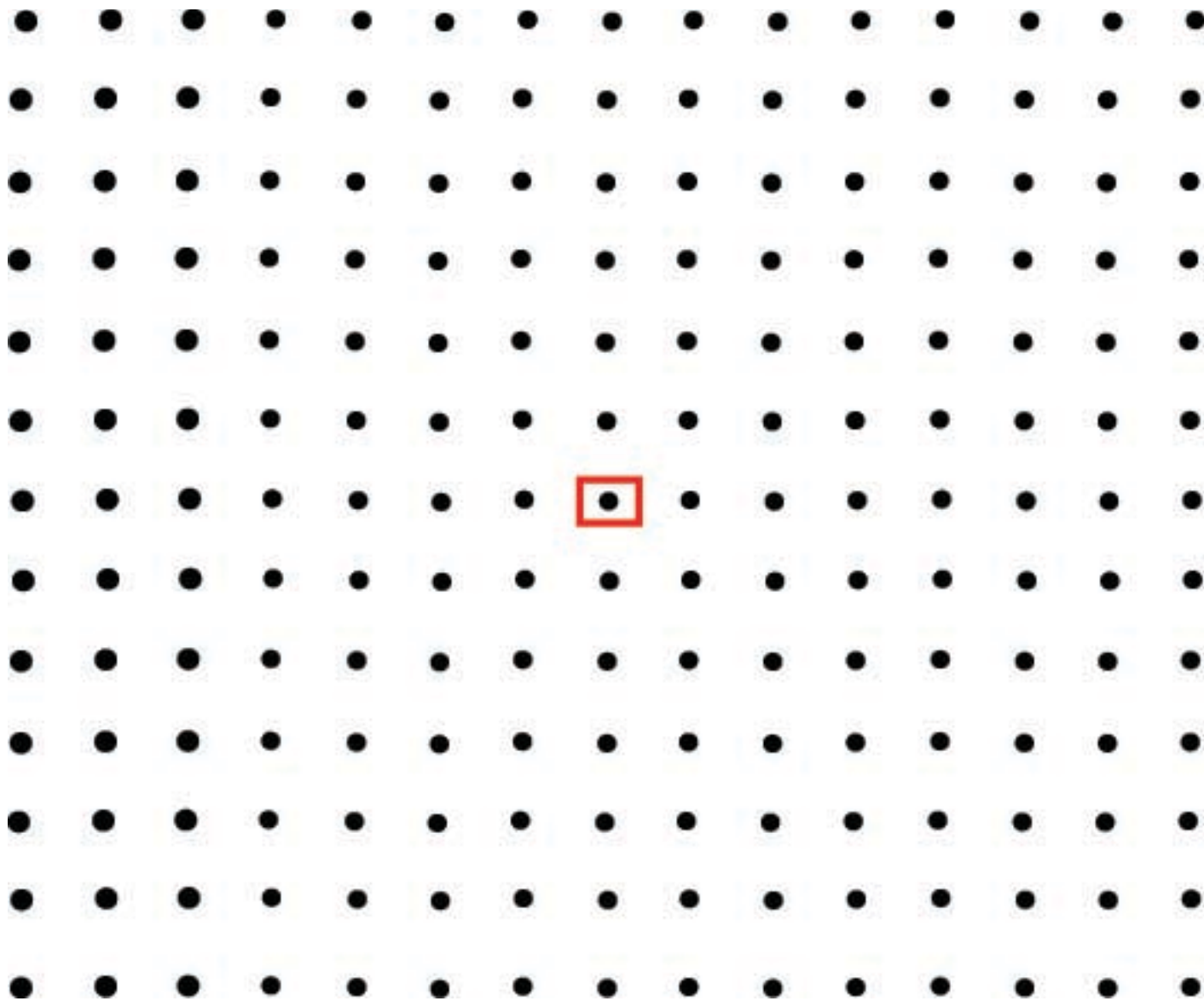


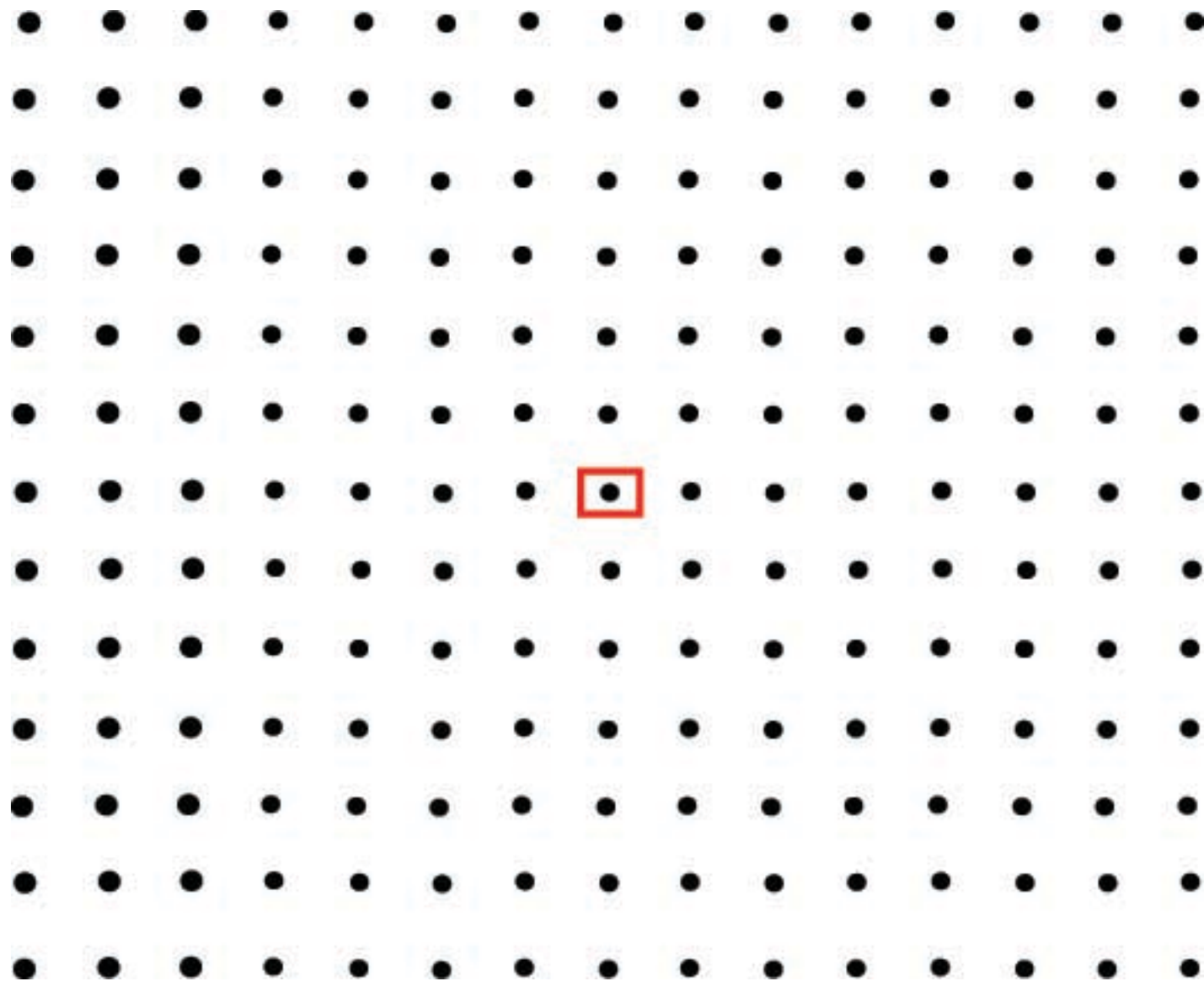


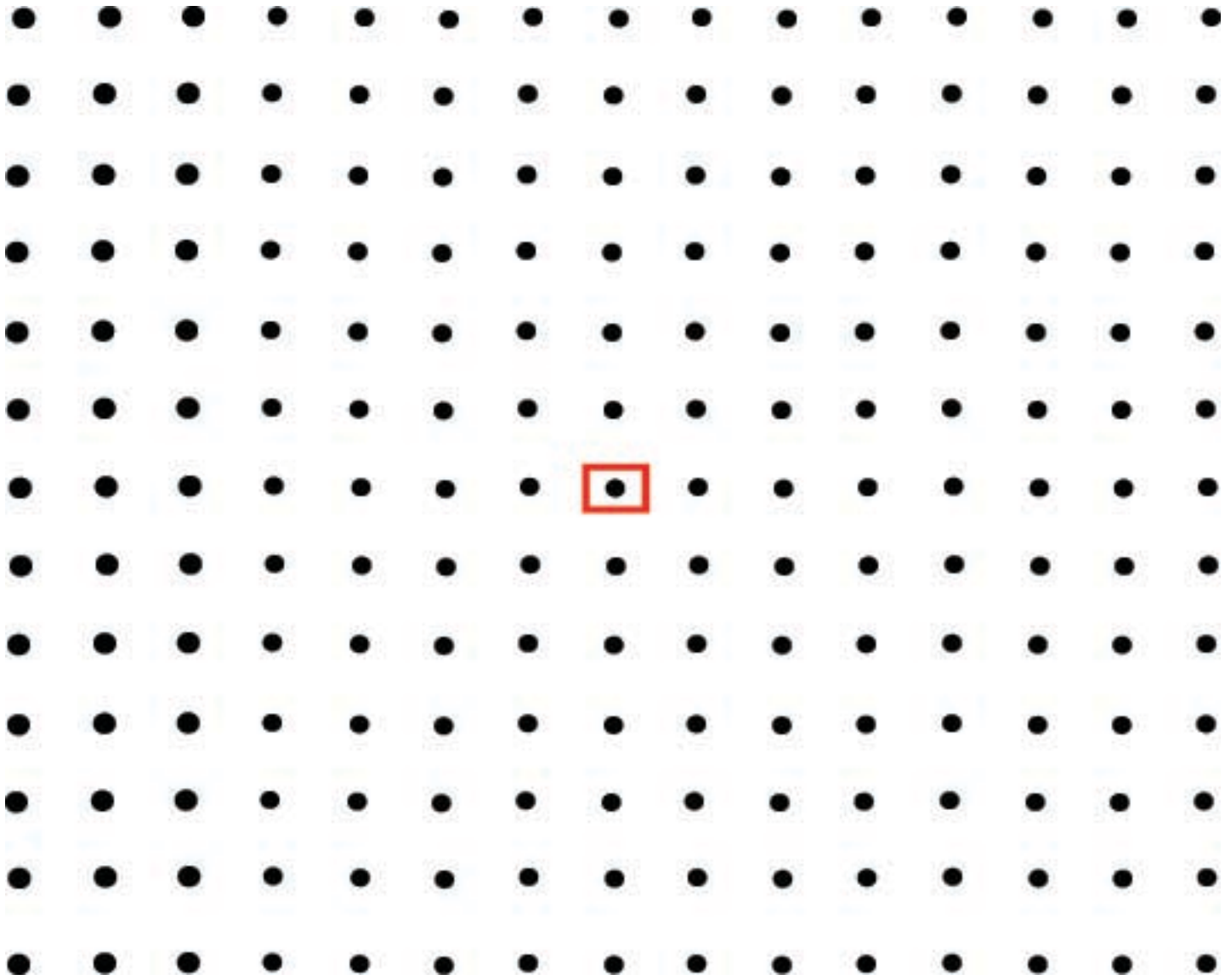


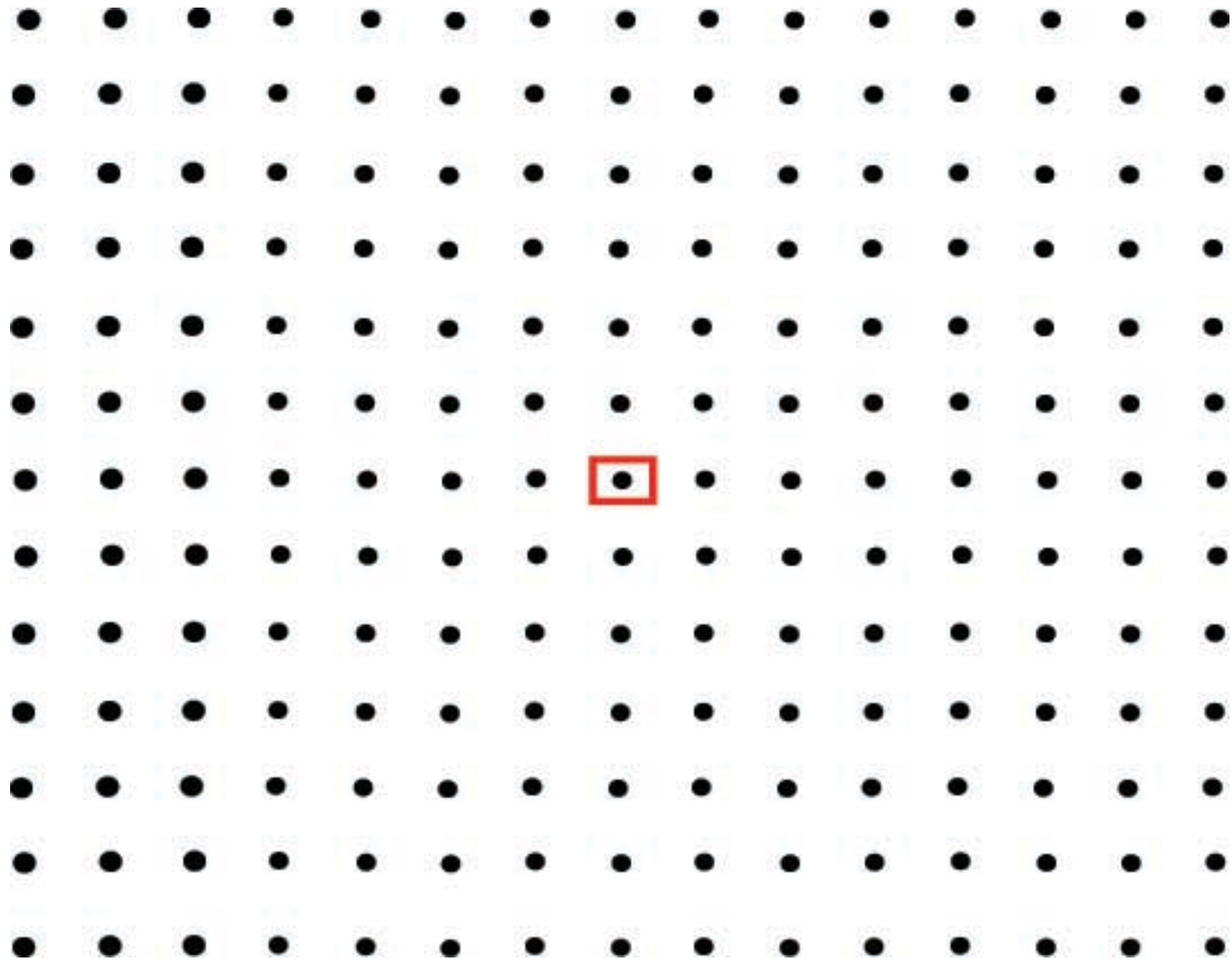


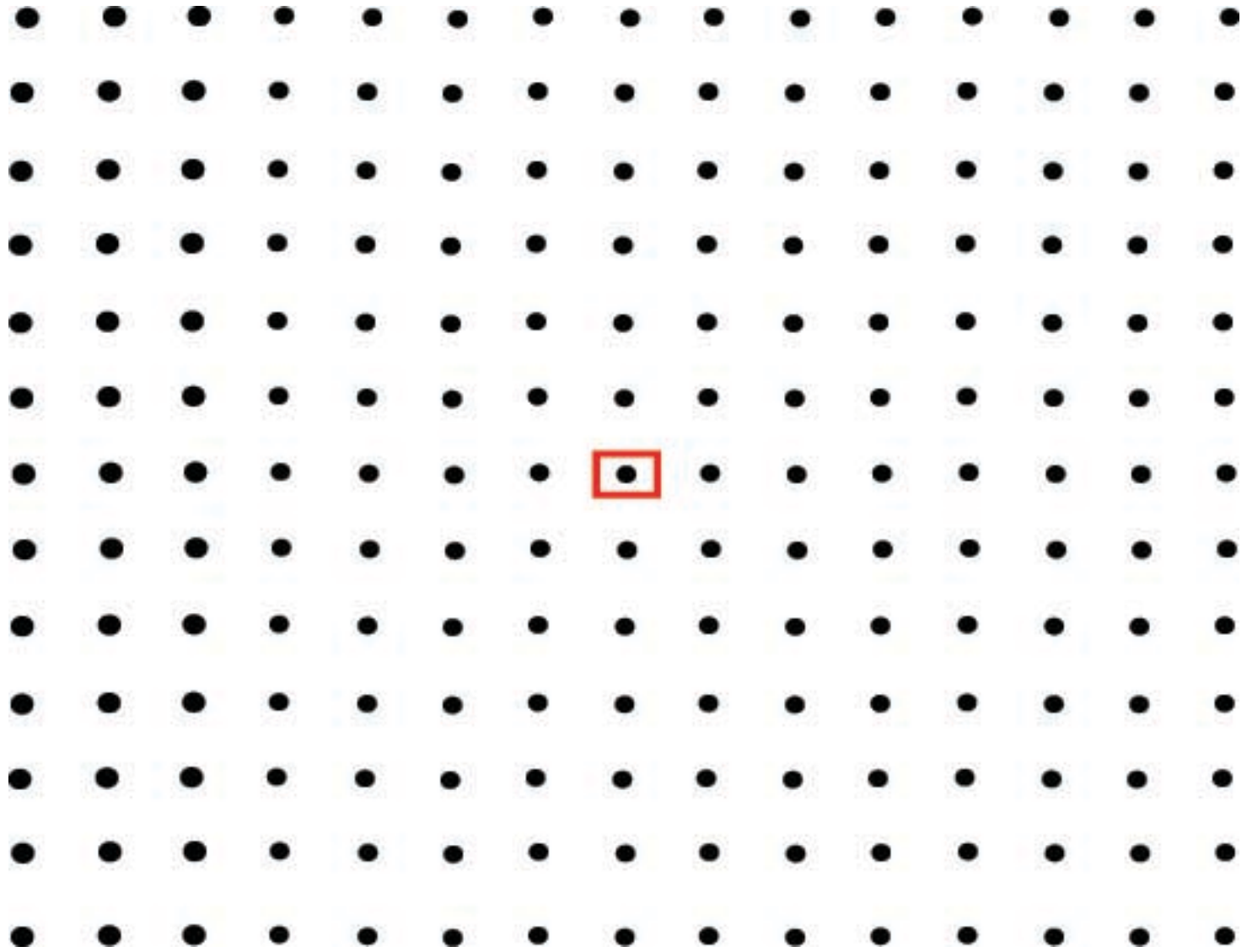


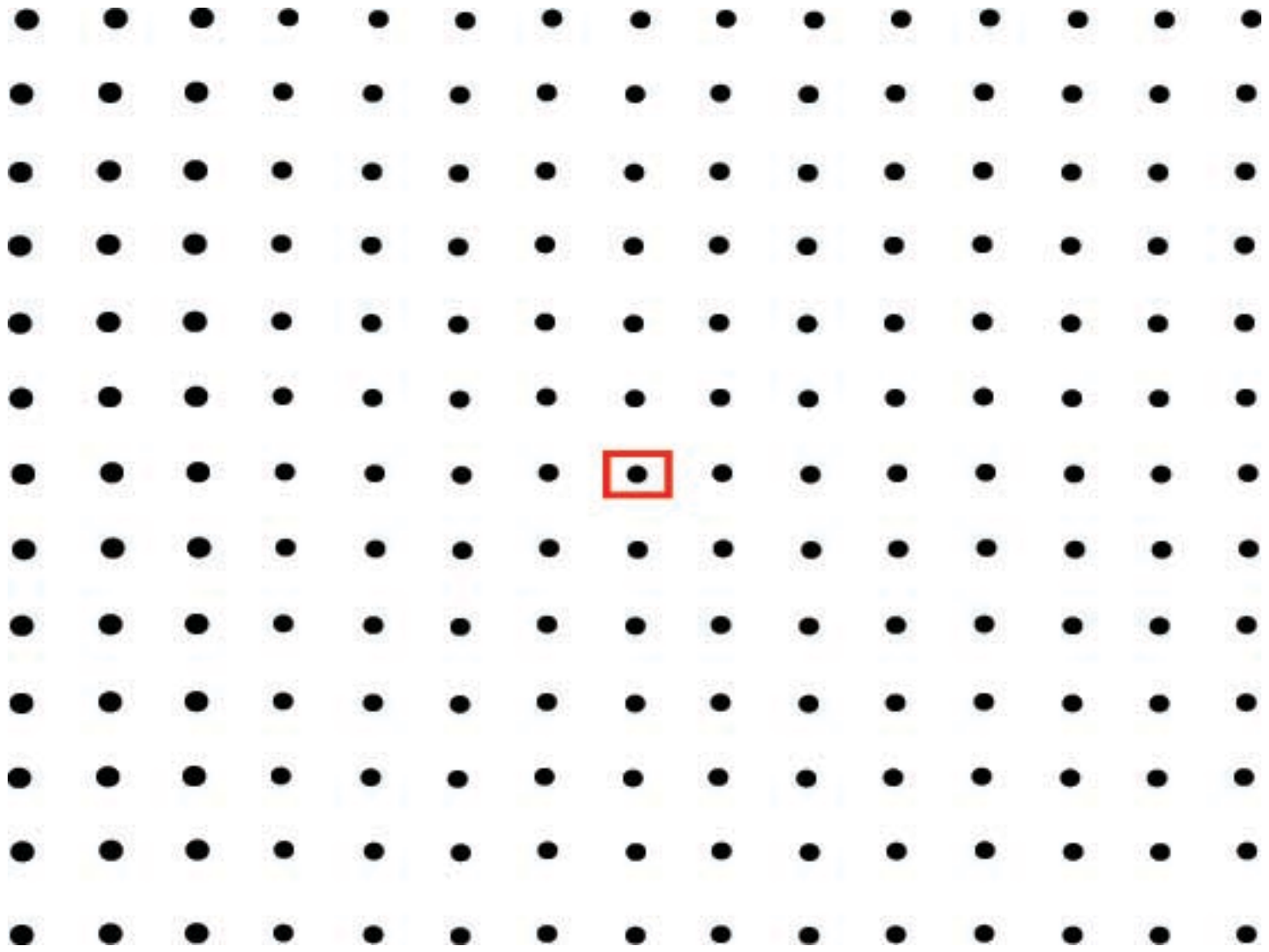


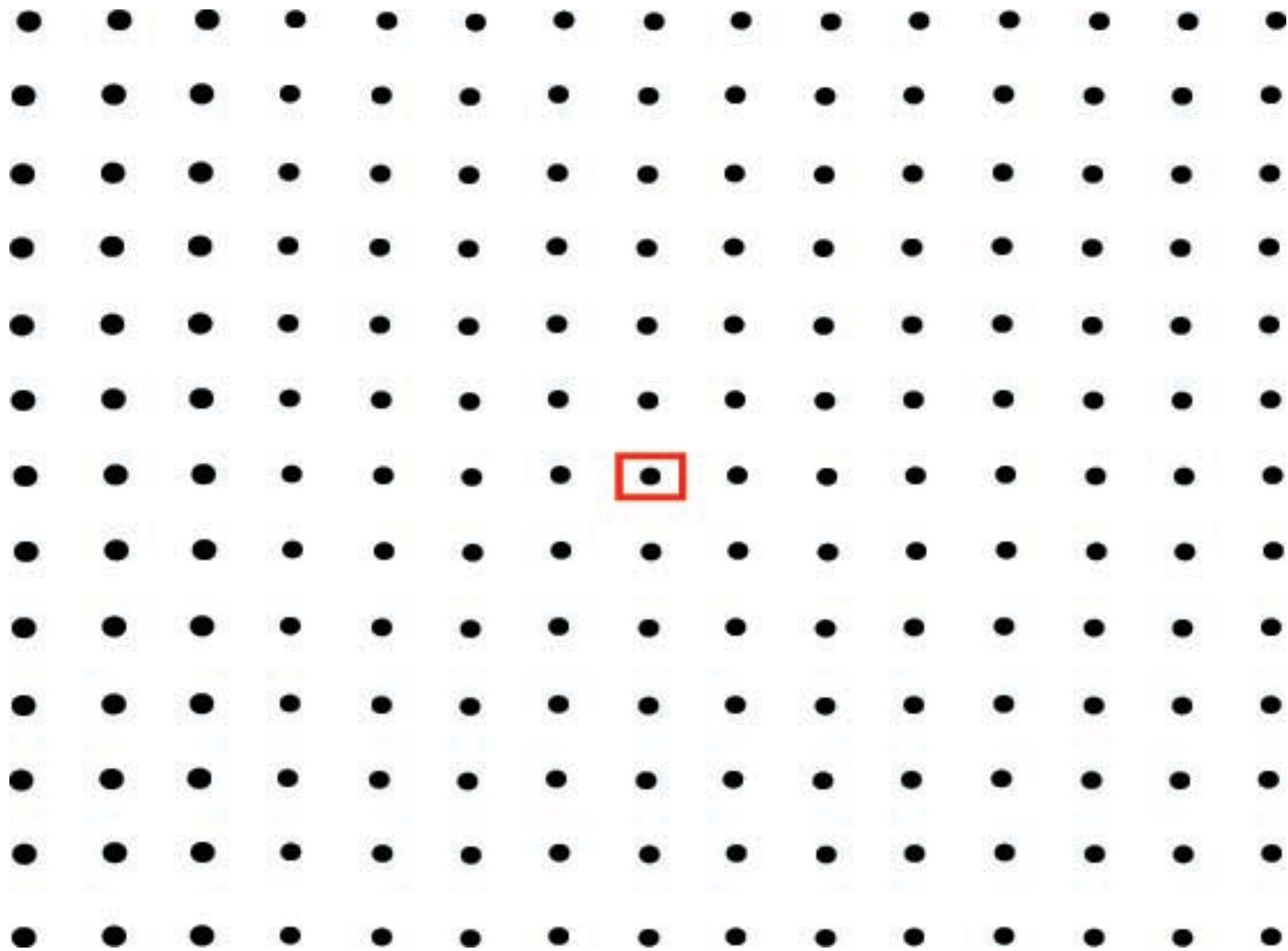


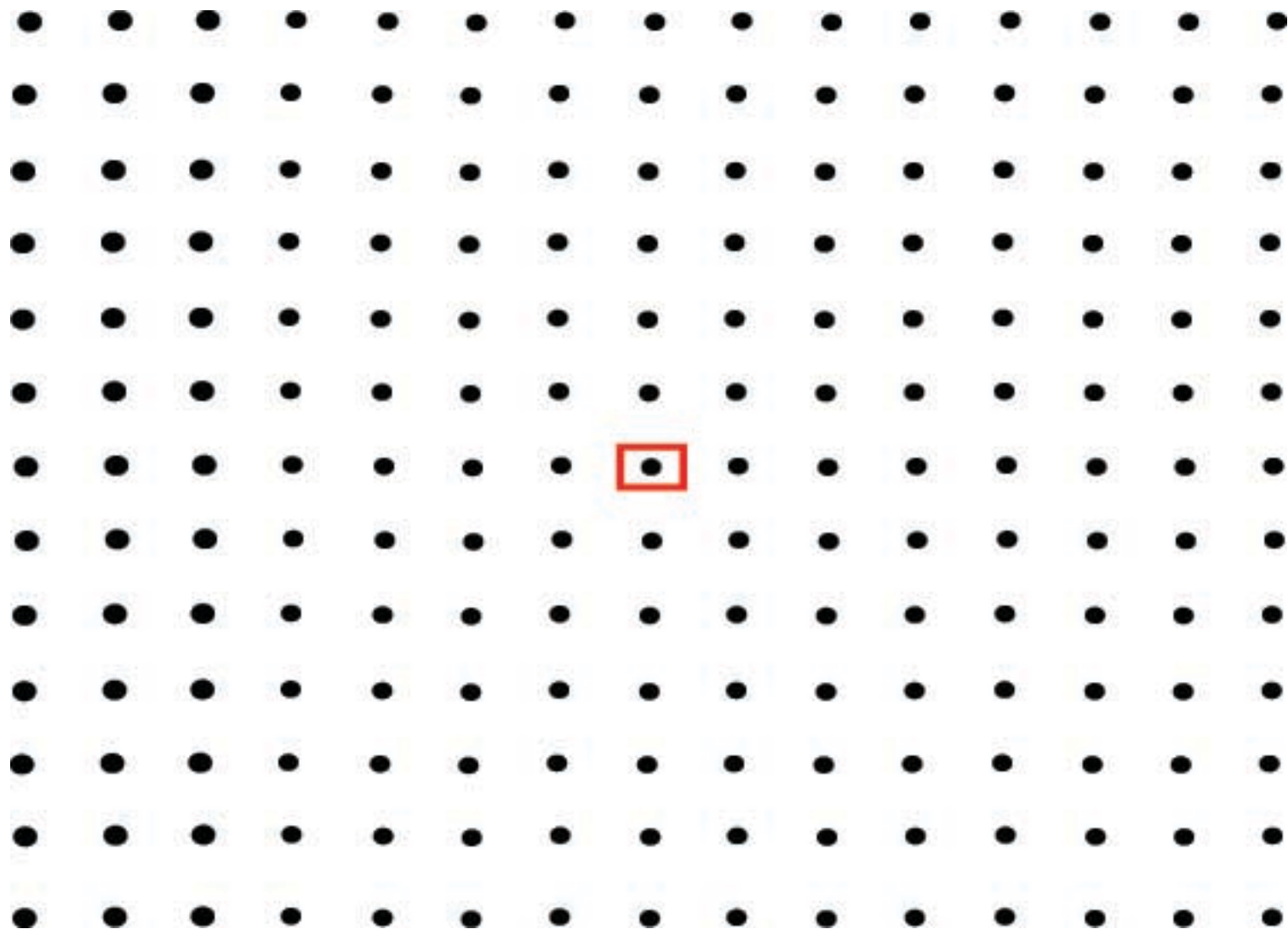


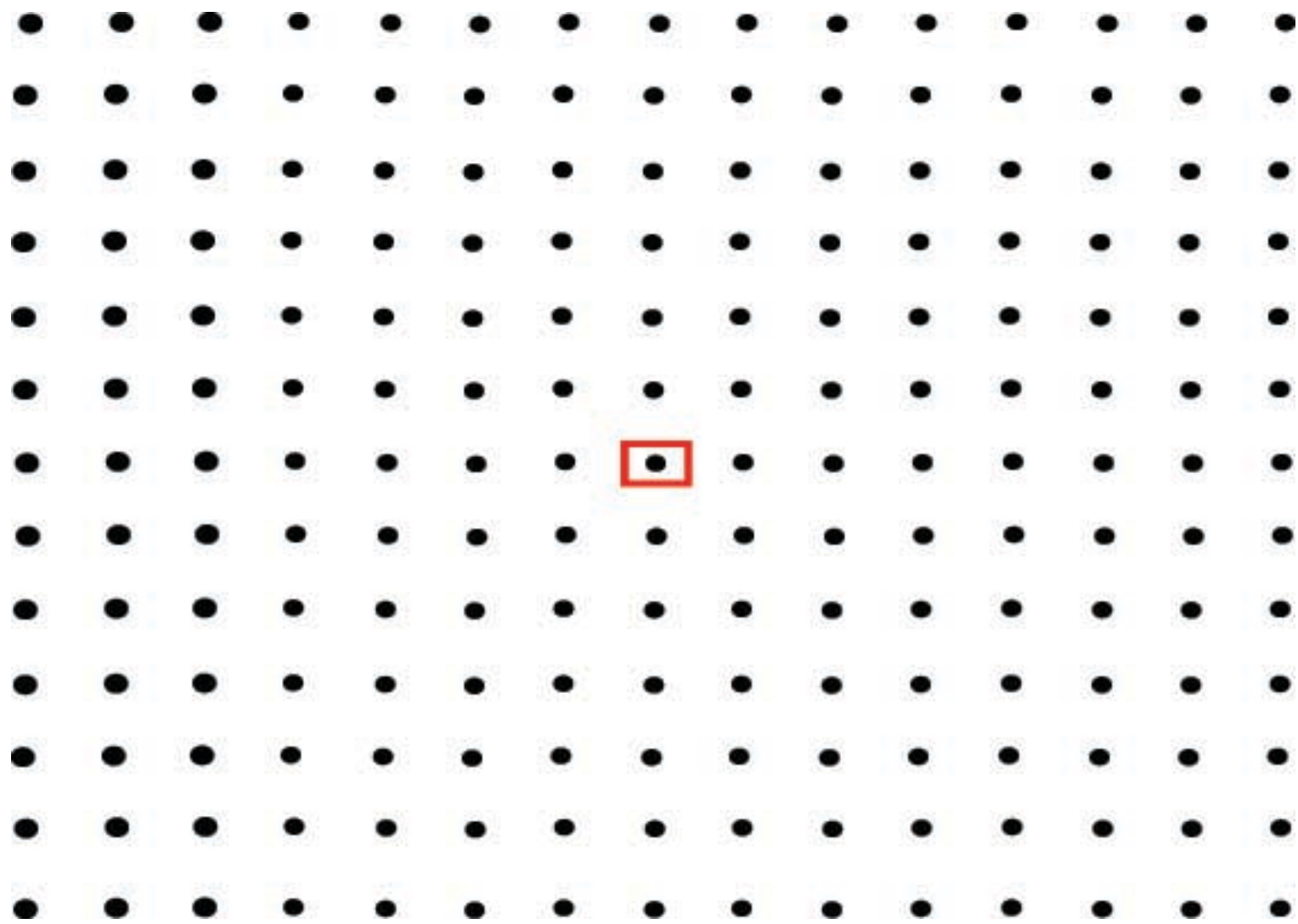


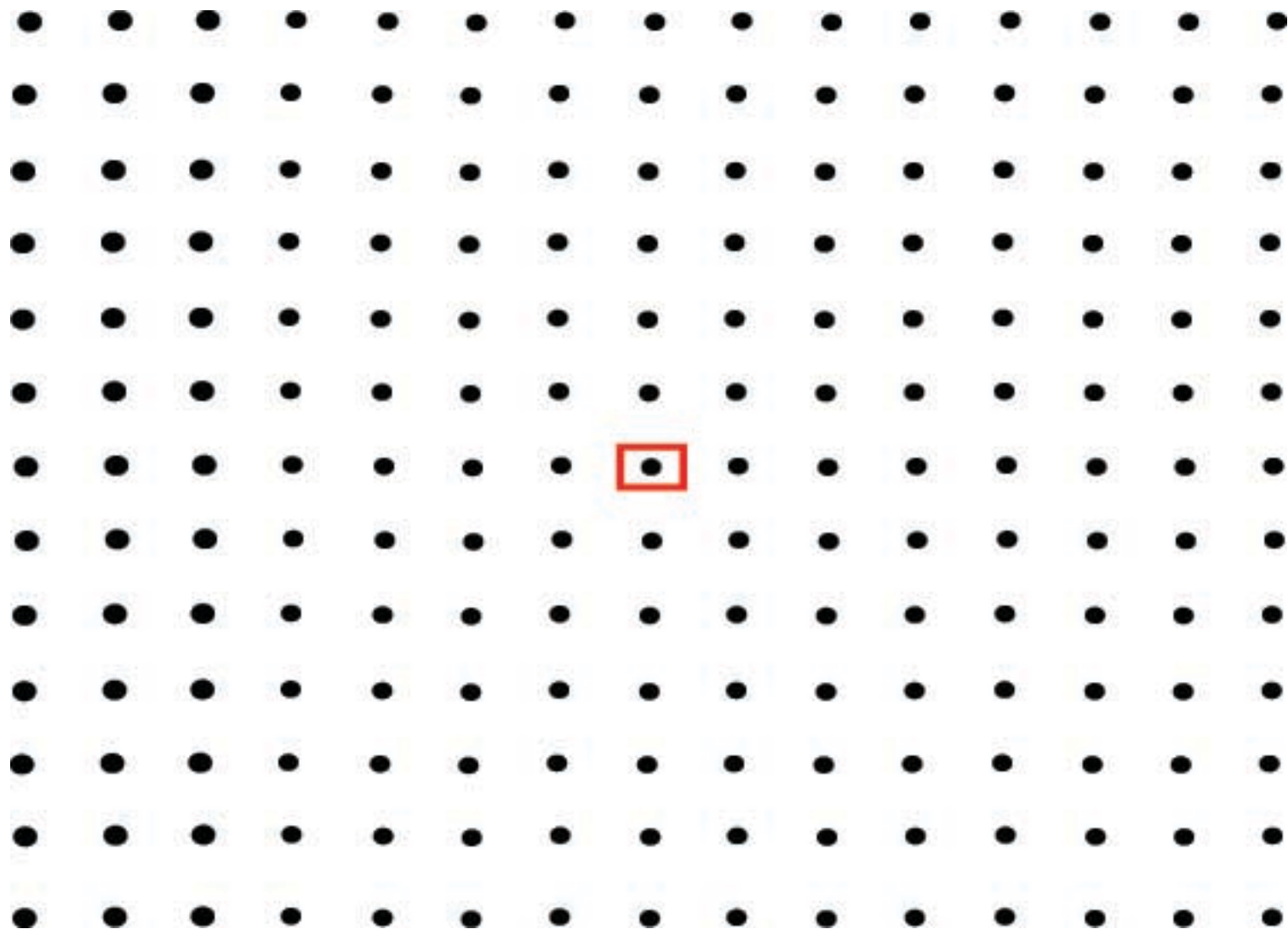


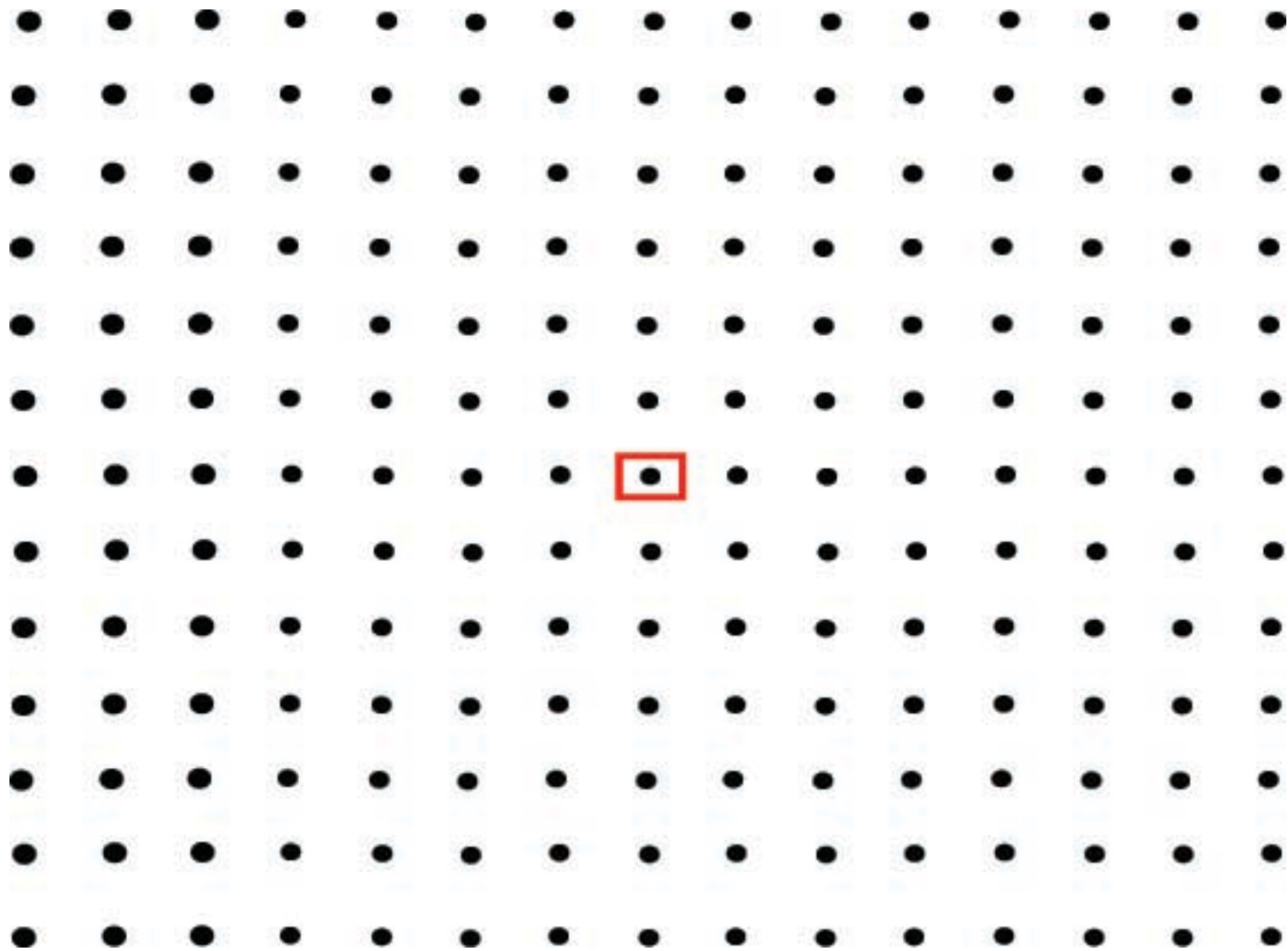


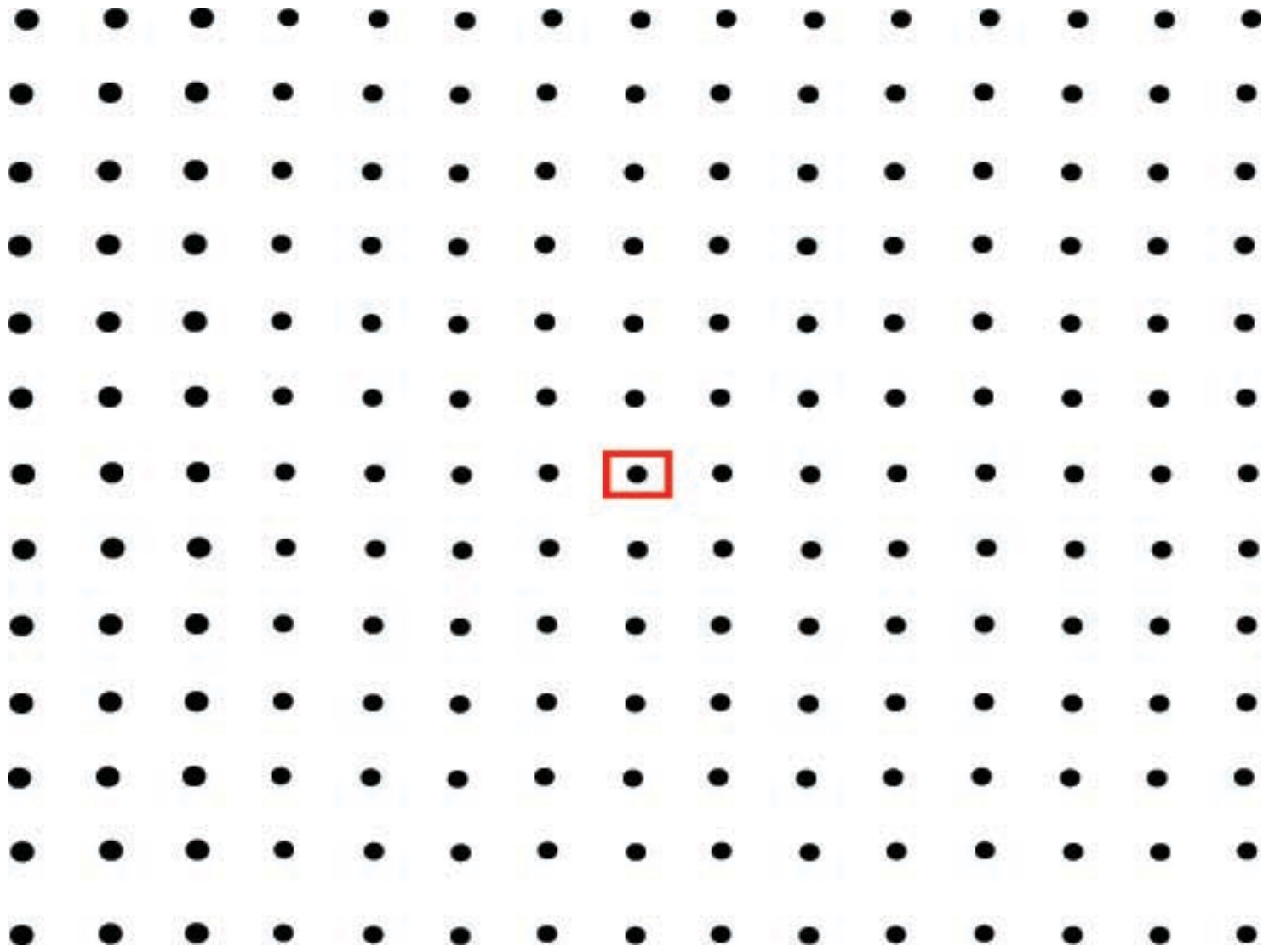


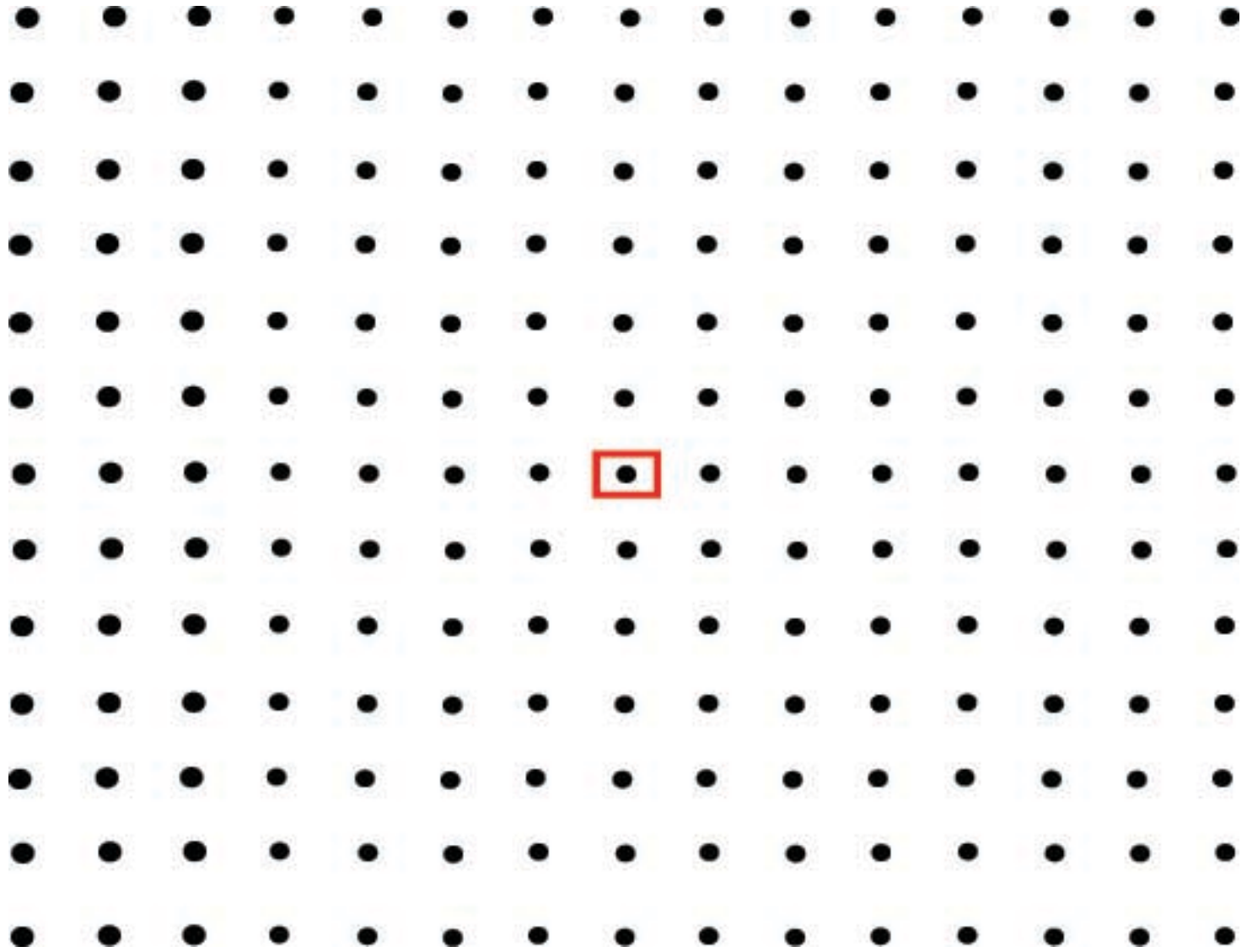


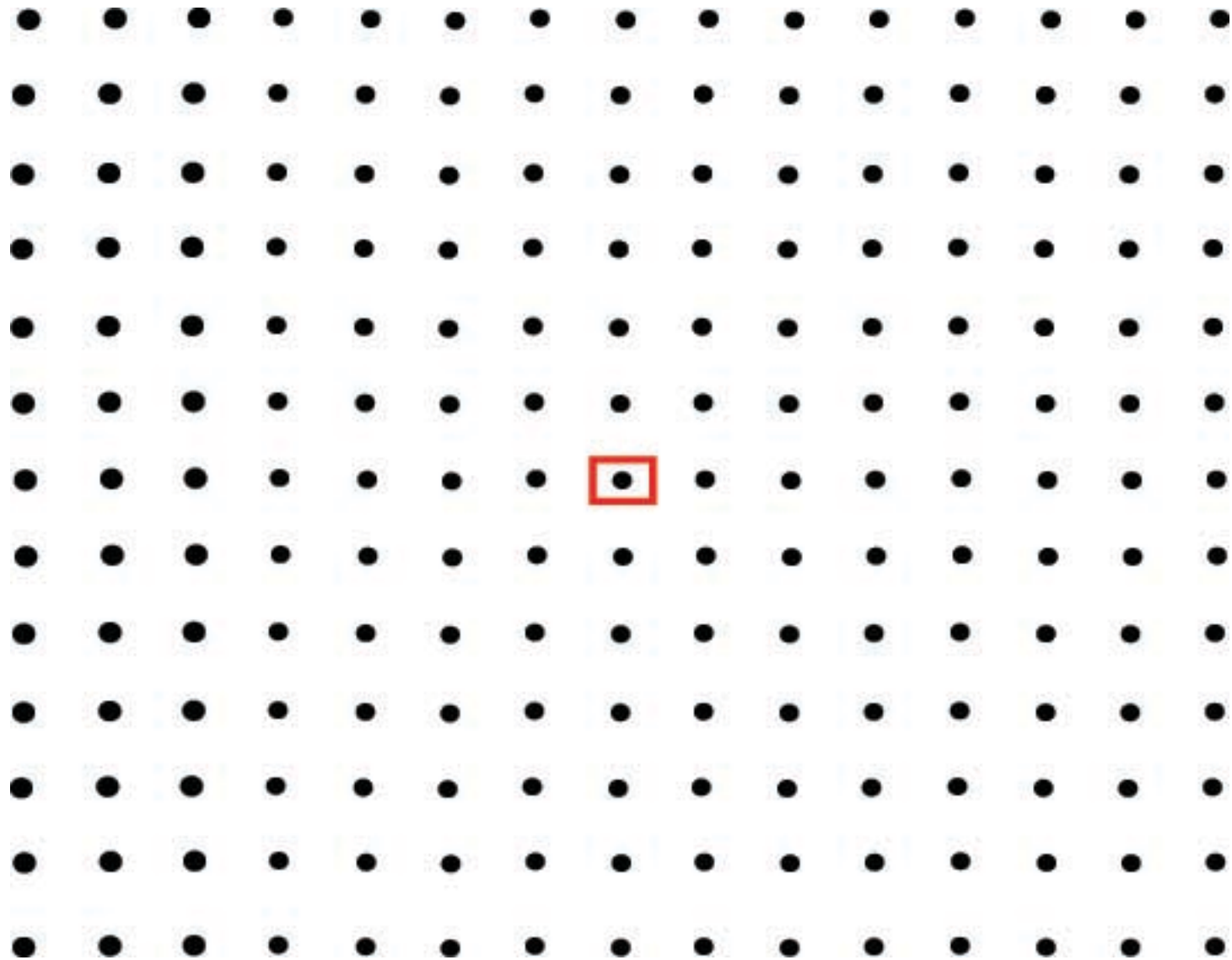


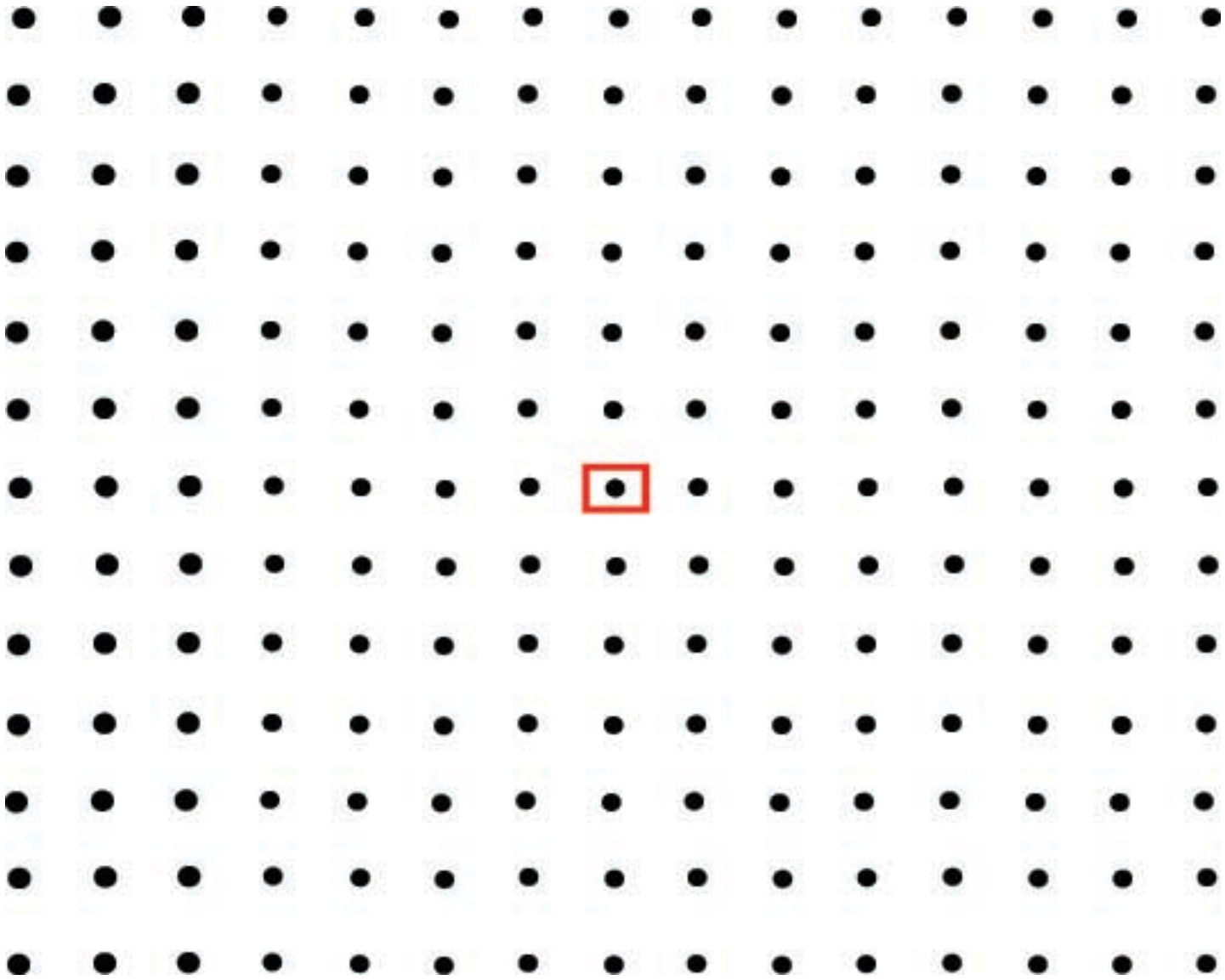


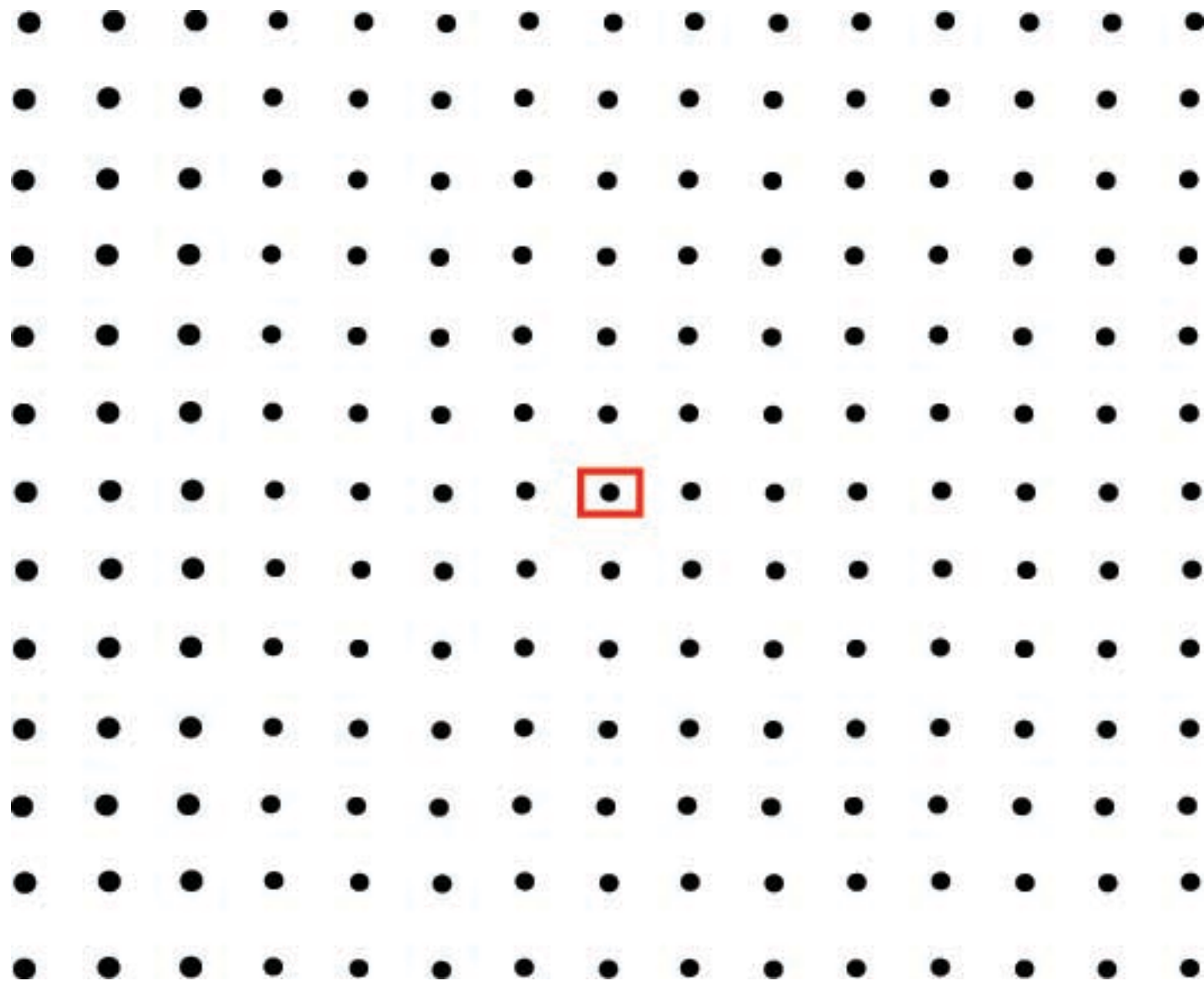


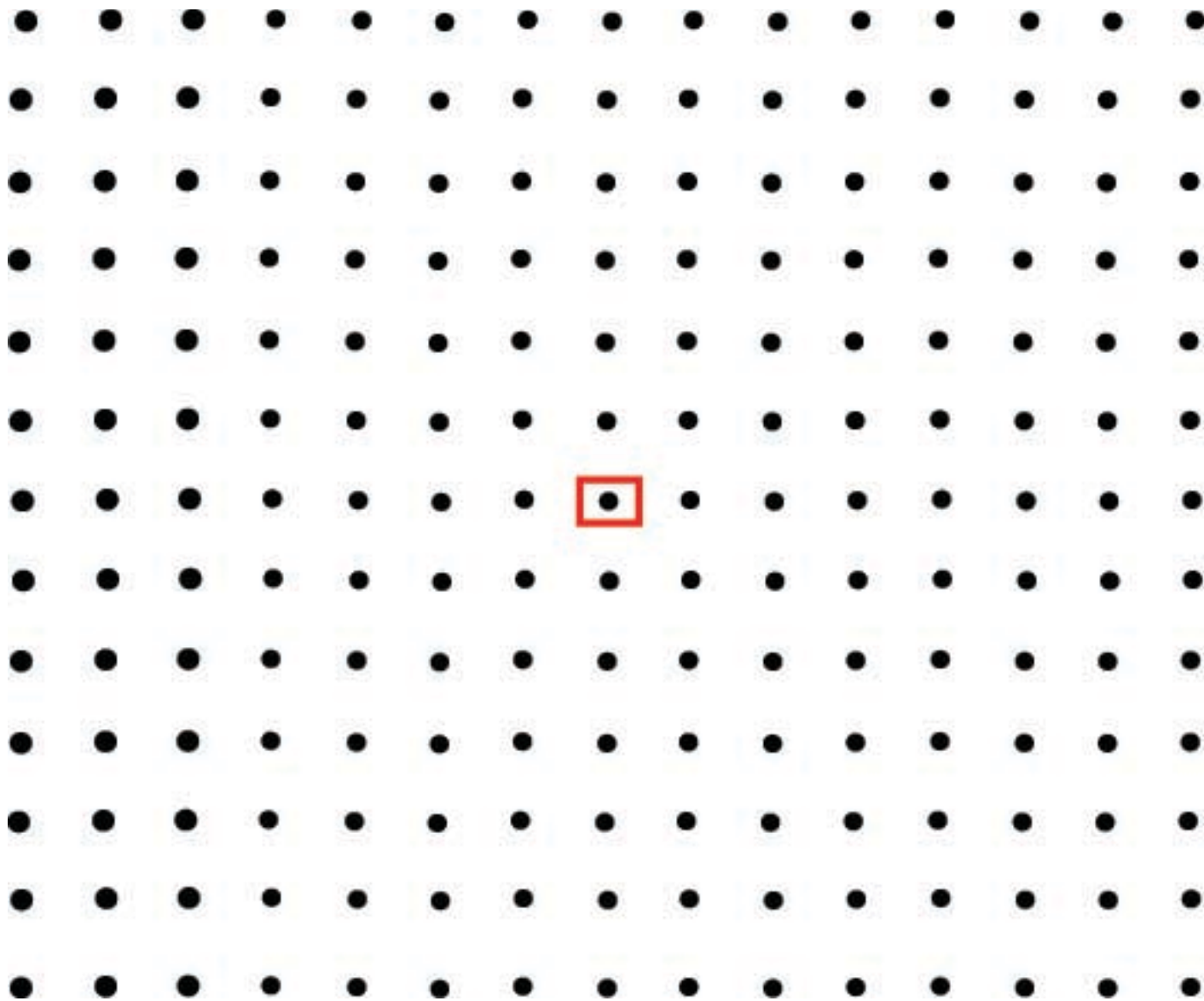


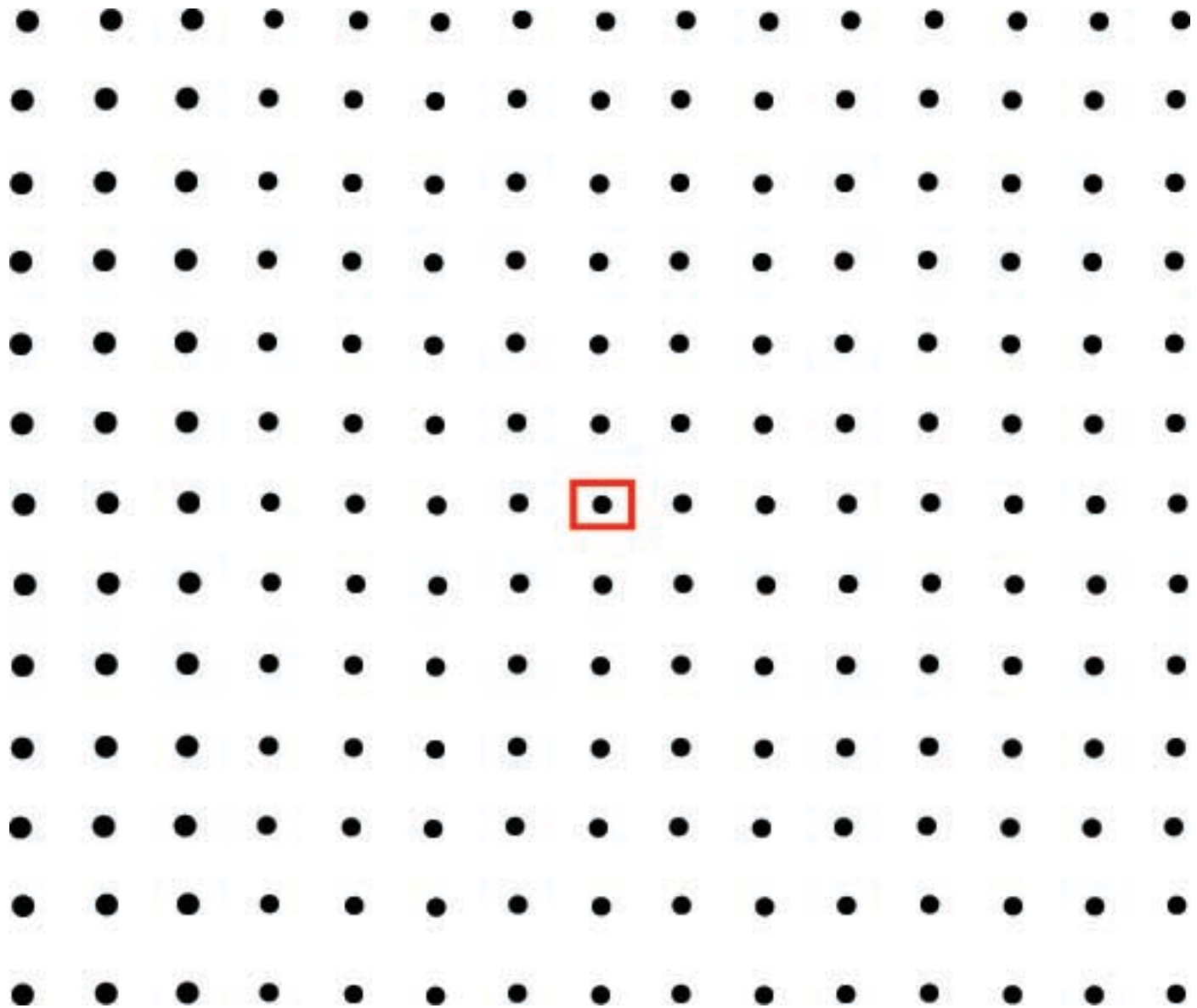












Relations for gravitational waves in modern notation

Intensity: $S_g = \frac{c^3}{16\pi G} \langle \dot{h}_+^2 + \dot{h}_x^2 \rangle$ pseudo tensor

$$\frac{c^3}{16\pi G} = 7.8 \times 10^{36} \text{ erg sec/cm}^2$$

Power radiated: $P_g = \frac{32 G m^2 x_0^4 \omega^6}{5 c^5}$ quadrupole formula

Relation to estimate GW amplitude: $h \approx \frac{\varphi_{\text{Newton}}}{c^2} \frac{v^2}{c^2} = \frac{Gm}{Rc^2} \frac{v^2}{c^2}$

1916 example: binary star system

$m_1 = m_2 = 1$ solar mass
 $T_{\text{orbit}} = 1$ day
 $R = 10$ Kly

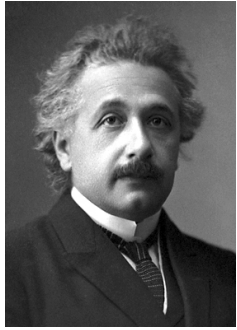
$h \sim 10^{-23}$ @ $\frac{1}{2}$ day period

$$Q = \frac{2\pi E_{\text{stored}}}{\Delta E_{\text{1period}}} \sim 10^{15} \quad \text{decaytime} \sim 10^{13} \text{ years}$$

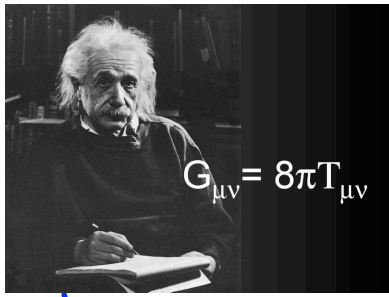
scaling $h \propto \frac{m^{\frac{5}{3}}}{RT_{\text{orbit}}^3}$

$$Q \propto \left(\frac{T_{\text{orbit}}}{m} \right)^{\frac{5}{3}}$$

— theory
— observation
— technology



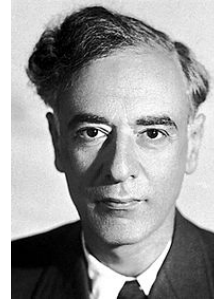
A. Einstein
Special Relativity
Random processes



A. Einstein
General Relativistic waves



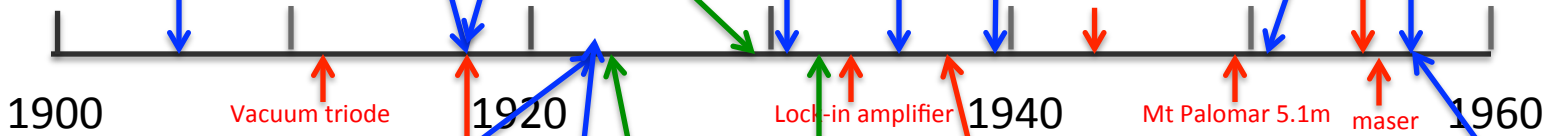
J.R. Oppenheimer
H. Snyder
Gravitational collapse to a BH



L. Landau & E. Lifshitz
Classical Theory of Fields



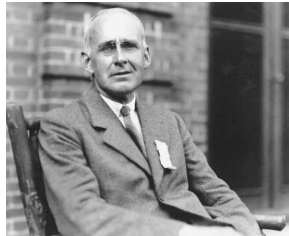
R.V. Pound
Cavity freq. stabilization



mostly mathematics

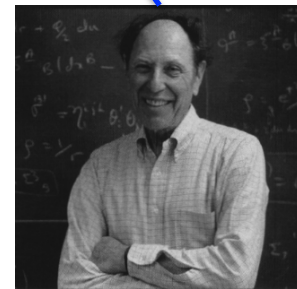
Feynman comments the field needs experiment, less mathematics

A.S. Eddington
skeptical: about pseudo tensor, inability to solve binary system, coordinate waves that propagate with the speed of thought also ones that might carry energy



M. Abraham
Electromagnetic analog

Understanding servoes
H. Bode
H. Nyquist
C. Shannon

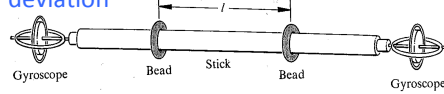


Josh Goldberg
US AirForce



F.A.E. Pirani, GW geodesic deviation

H. Bondi : beads on rod, energy in GW



J. Weber Acoustic resonant bar GW detectors



V. Braginsky (1974) quantum limit



P. Bender LISA (1975)



J. Taylor & R. Hulse binary pulsar



D.C. Backer

Space program: C. Misner, R.V. Pound, P. Bender, RW Laser GW detector, X and K band ranging

GW from inflation

millisecond pulsars

Millisecond pulsar timing GW detection

Chapel Hill meeting J. Goldberg

M. Gertsenshtein GW interaction with EM waves Kerr metric

Integrated circuit transistor

Kruskal coordinates Laser

X-ray astronomy

J. Wheeler: Physical reality of GW, Weber bar

1960



R.H. Dicke: Eotvos experiment, electronic cooling of mechanical instrument



J. Bell & A. Hewish pulsars

CMB discovery

BH in Cygnus X-1

PDP 11 Lab computer

Electromagnetically coupled GW detector (RW)

F. Estabrook H. Wahlquist space craft ranging

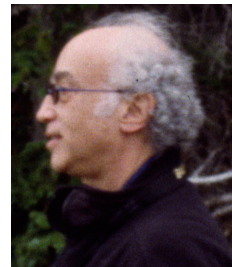


A. Guth

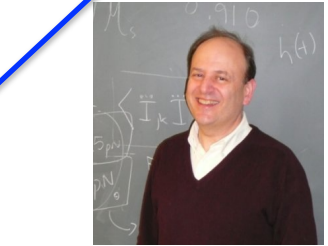


A. Linde

Battelle Seattle workshop



R. Isaacson



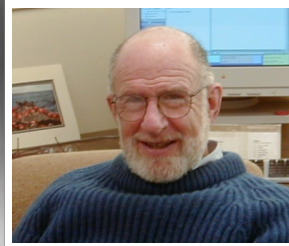
P. Saulson

1980

NSF report 1983 LIGO



B. McDaniel



A. Sessler

1990

COBE cosmic structure

R. Vogt becomes 1st Director of LIGO

NSF Study Committee on Interferometric GW detection



Acoustic bar GW Detector groups



R. Garwin



W. Fairbank



E. Amaldi

1965-1975
Room T bars

Bell Labs
Frascati
Glasgow
IBM
Rochester
Max Planck
Rome

1975-1990+
Cryogenic bars

Frascati
Louisiana
Moscow
Perth
Rochester
Stanford

2000 ->
Spherical cryogenic detectors

Brazil
Netherlands



A. Tyson



W. Hamilton

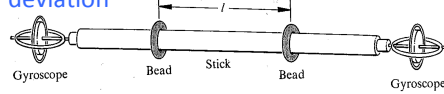


P. Michelson



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Cosmic inflation

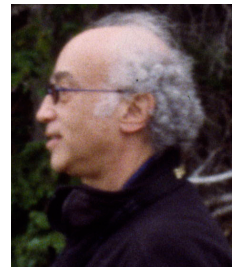


A. Guth

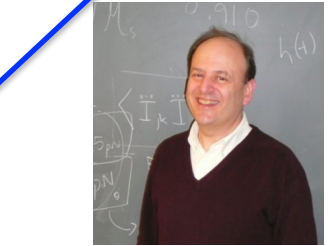


A. Linde

Battelle Seattle workshop



R. Isaacson



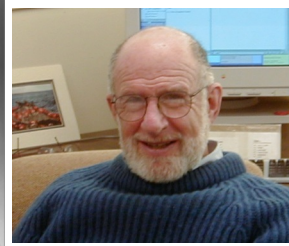
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NSF Study Committee on Interferometric GW detection



Plane gravitational waves

Transverse Plane Wave Solutions with “Electric”
and “Magnetic” Terms

Geometric Interpretation

$$ds^2 = g_{ij} dx^i dx^j$$

$$g_{ij} = \eta_{ij} + h_{ij} \quad \text{weak field}$$

$$\eta_{ij} = \begin{pmatrix} 1 & & & \\ & -1 & & \\ & & -1 & \\ & & & -1 \end{pmatrix} \quad \begin{array}{l} \text{Minkowski Metric of} \\ \text{Special Relativity} \end{array}$$

Gravity Wave Propagating in the x_1 Direction

$$h_{ij} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & h_{22} & h_{23} \\ 0 & 0 & h_{32} & h_{33} \end{pmatrix} \quad \text{all } h_{ij} \ll 1$$

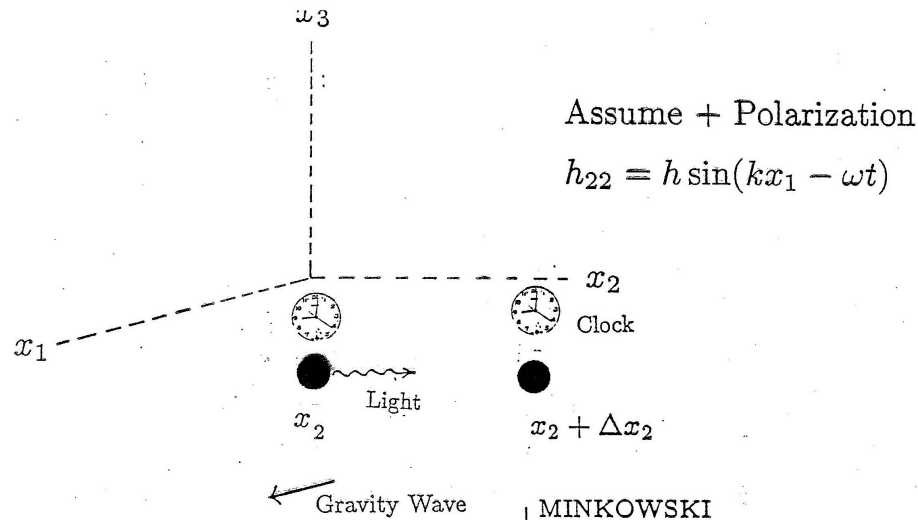
Plane Wave

$$\mathbf{h}_{22} = -\mathbf{h}_{33} \quad \mathbf{h}_{23} = \mathbf{h}_{32}$$

+ polarization × polarization

And All Only Function of $x_1 - ct$

Timing light in the gravitational wave



$$\Delta s^2 = 0 = c^2 \Delta t^2 - \left(1 + h \sin(kx_1 - \omega t)\right) \Delta x_2^2$$

LIGHT RAY

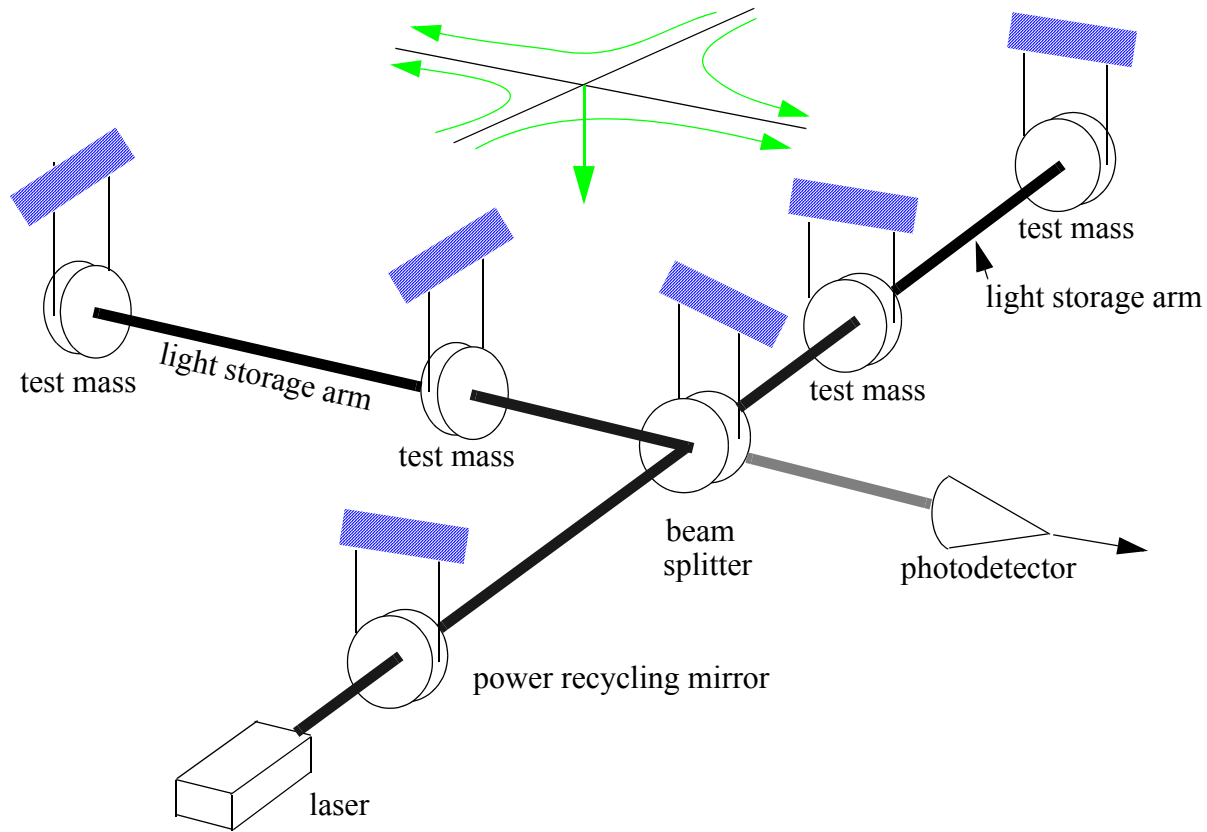
Let $\Delta t \ll \frac{1}{\omega}$ $h \ll 1$

$$c \Delta t \cong \left(1 + \frac{h}{2} \sin(kx_1 - \omega t)\right) \Delta x_2$$

←
 INFERRED
 DISTANCE
 BETWEEN POINTS

$$\frac{\delta(c \Delta t)}{\Delta x_2} = \frac{h}{2} \sin(kx_1 - \omega t) \quad \text{Time Dependent Strain}$$

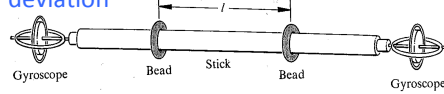
$$\frac{\Delta l}{l} = \frac{h}{2} \quad \text{The Measurable Quantity}$$





F.A.E. Pirani, GW geodesic deviation

H. Bondi : beads on rod, energy in GW



J. Weber Acoustic resonant bar GW detectors



V. Braginsky (1974) quantum limit



P. Bender LISA (1975)



J. Taylor & R. Hulse binary pulsar



D.C. Backer

Space program: C. Misner, R.V. Pound, P. Bender, RW Laser GW detector, X and K band ranging

GW from inflation

millisecond pulsars

Millisecond pulsar timing GW detection

Chapel Hill meeting J. Goldberg

M. Gertsenshtein GW interaction with EM waves Kerr metric

Integrated circuit transistor

Kruskal coordinates Laser

X-ray astronomy

J. Wheeler: Physical reality of GW, Weber bar

1960

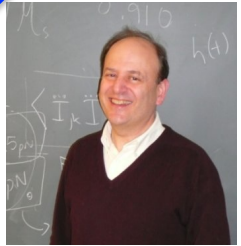
J. Bell & A. Hewish pulsars CMB discovery

1970 Electromagnetically coupled GW detector (RW)

BH in Cygnus X-1 PDP 11 Lab computer

Battelle Seattle workshop F. Estabrook H. Wahlquist space craft ranging

1980 NSF report 1983 LIGO



P. Saulson

1990 COBE cosmic structure

R. Vogt becomes 1st Director of LIGO

NSF Study Committee on Interferometric GW detection



R.H. Dicke: Eotvos experiment, electronic cooling of mechanical instrument



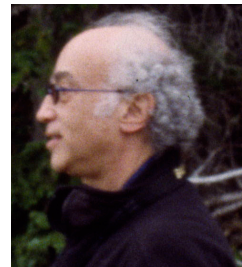
A. Guth



A. Linde



R. Isaacson



B. McDaniel

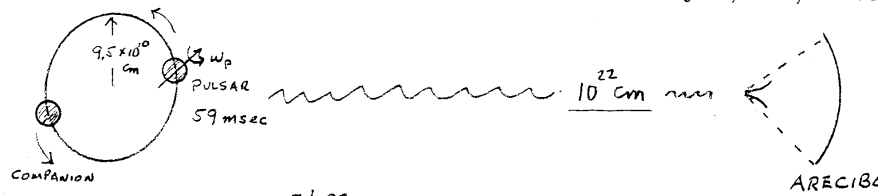


A. Sessler

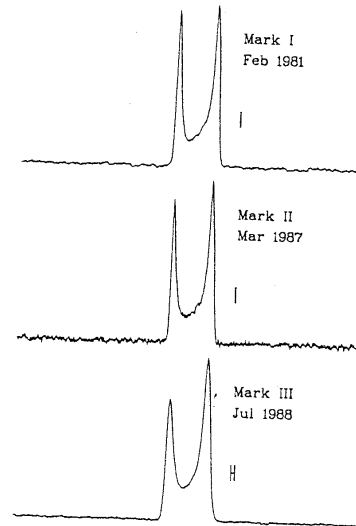
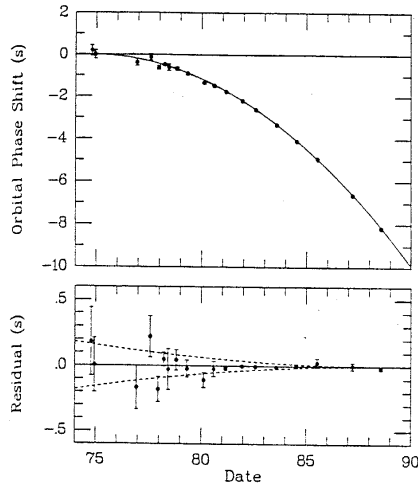


The Binary Pulsar PSR 1913 + 16

Taylor, J.H., Weisberg, J.M.



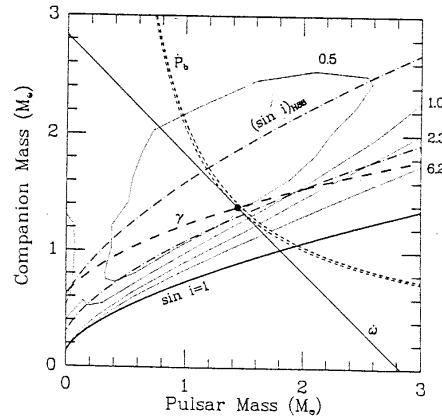
ORBITAL PERIOD $\sim 7.75 \text{ hrs}$
 $e \sim 0.6$ $\frac{v_{\text{ORB}}}{c} \sim 10^{-3}$ $\frac{GM}{2ac^2} \sim 10^{-6}$

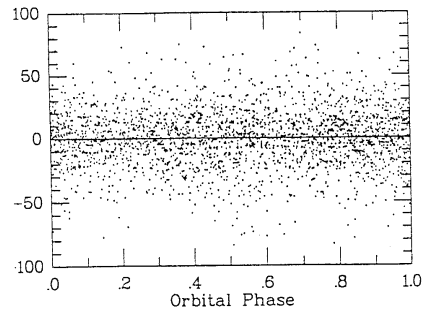


Detected Pulsar pulses

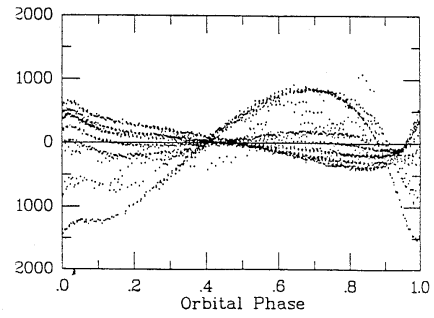
Energy loss to gravitational radiation

$$P = G(\ddot{Q})^2 / 45c^5$$





Post fit residuals



Post fit residuals if $\gamma \rightarrow 0$

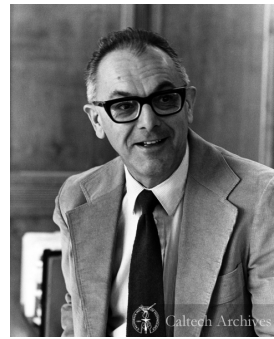
Taylor, J.H., Weisberg, J.M.



R. Drever



K. Thorne



R. Vogt

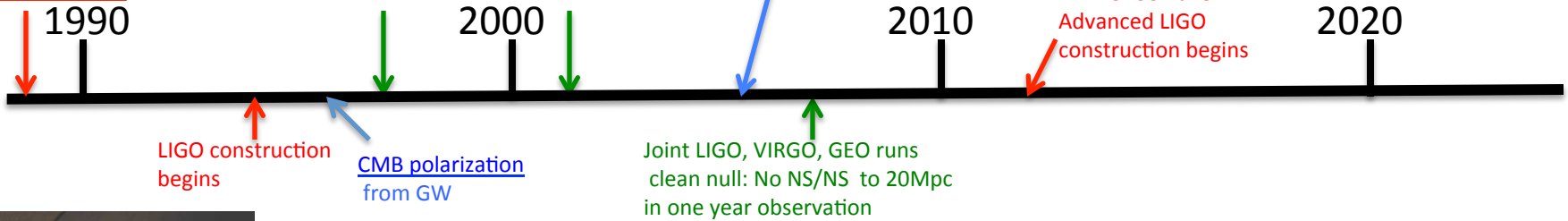


F. Pretorius
Numerical relativity waveforms



D. H. Shoemaker
Advanced LIGO
construction begins

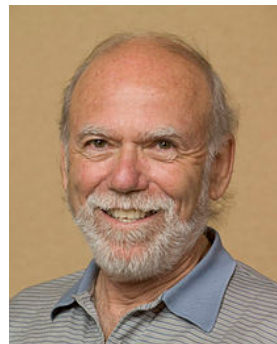
LIGO 1989
Proposal to NSF



A. Brillet & A. Giazotto
VIRGO 1990



K. Danzmann
GEO 1995 & LISA



B. Barish 2nd
LIGO director

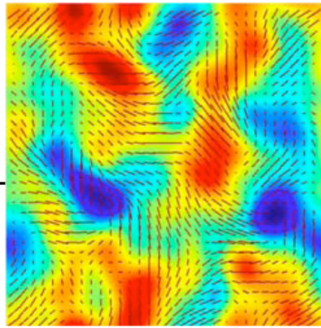


J. Marx 3rd
LIGO Director



D. Reitze 4th
LIGO Director

*Cosmic Microwave Background
Polarization B Modes*



h

10^{-5}

10^{-10}

10^{-15}

10^{-20}

10^{-25}

Primeval gravitational waves from inflationary epoch

Measured at epoch of recombination $z \sim 1000$ and reionization $z \sim 6$

Gravitational Wave Spectrum

Pulsar Timing



Supermassive BH coalescences

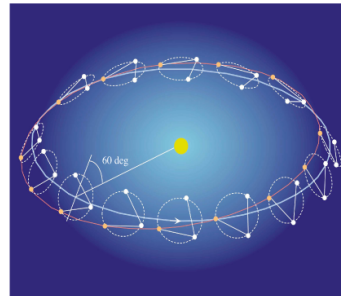
Isotropic GW background from unresolved sources

Massive BH coalescences

Small mass/BH infalls

White dwarf binaries in our galaxy

Space-based Interferometers



Compact binary coalescences: neutron stars and black holes

Asymmetric pulsar rotations

Ground-based Interferometers



10^{-16}

10^{-12}

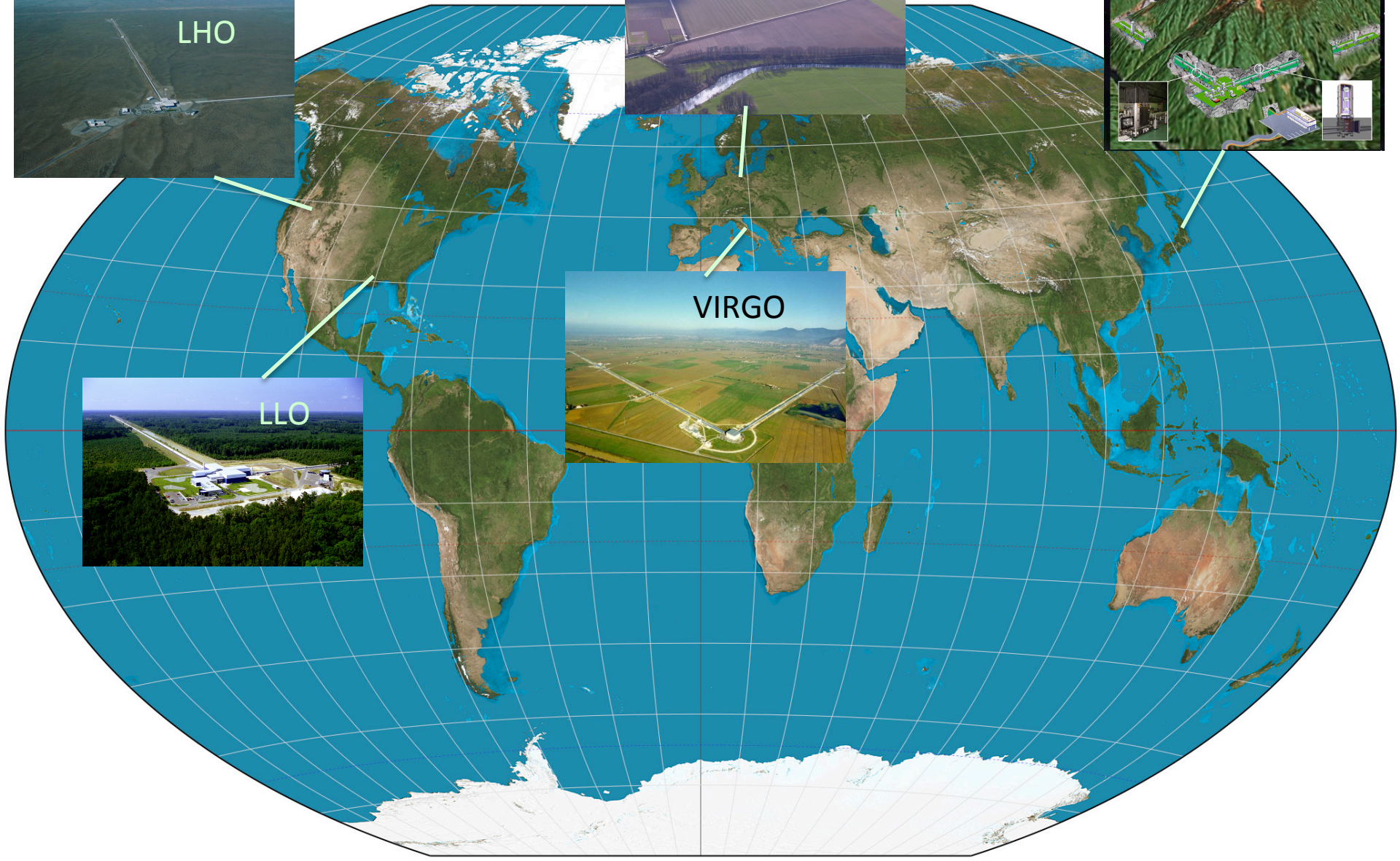
10^{-8}

10^{-4}

10^0

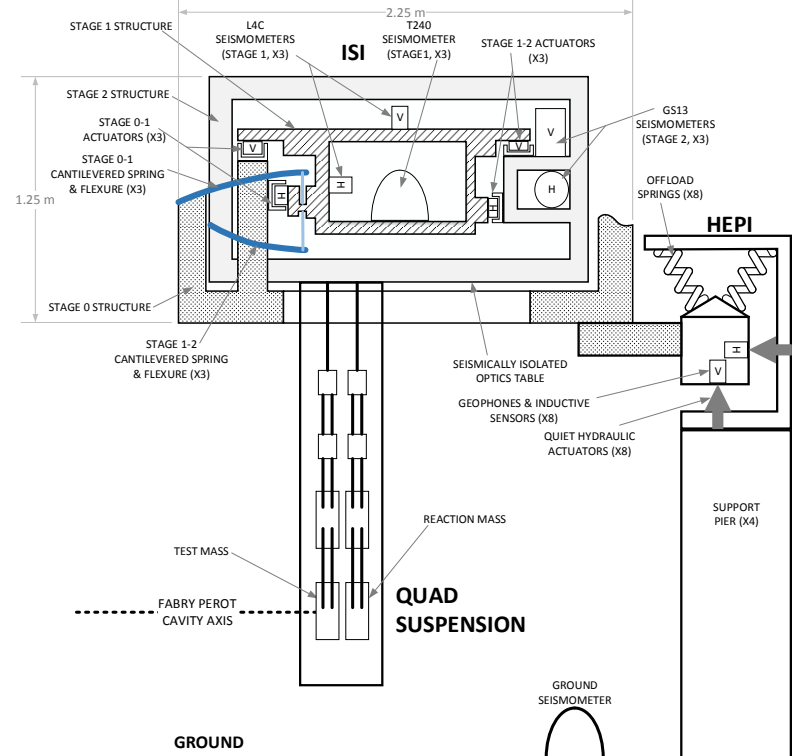
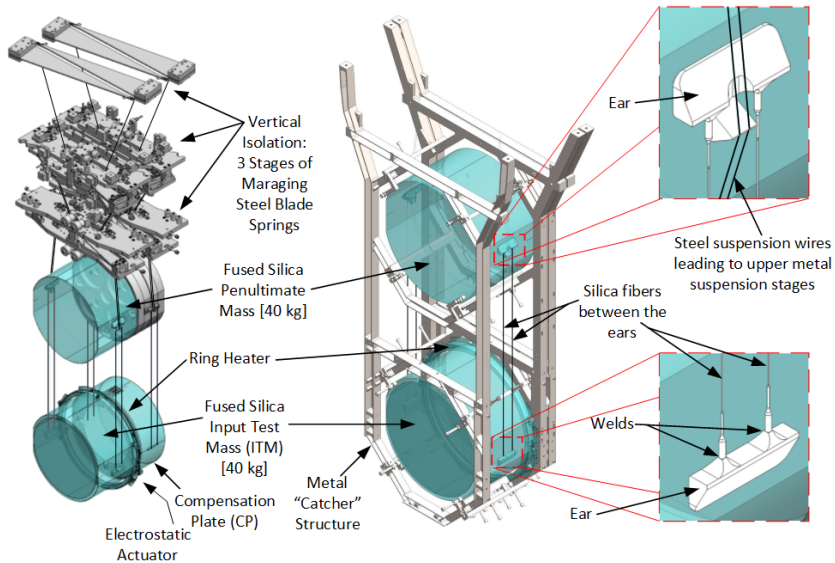
10^4

Frequency Hz

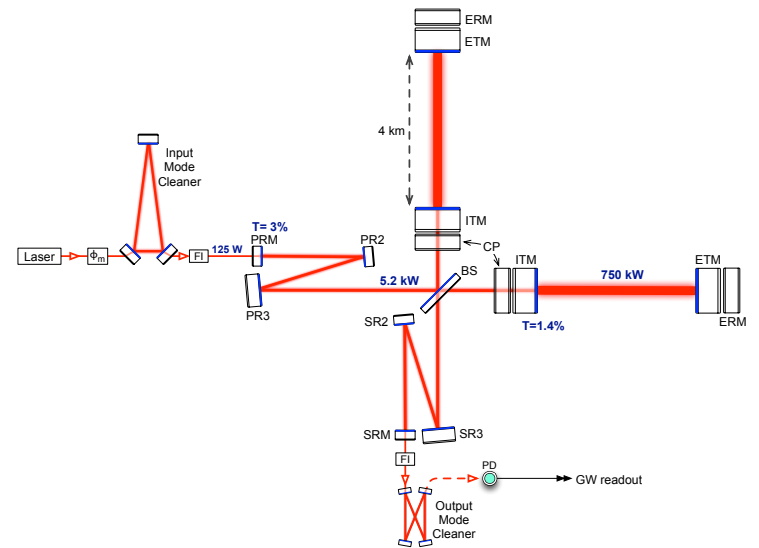
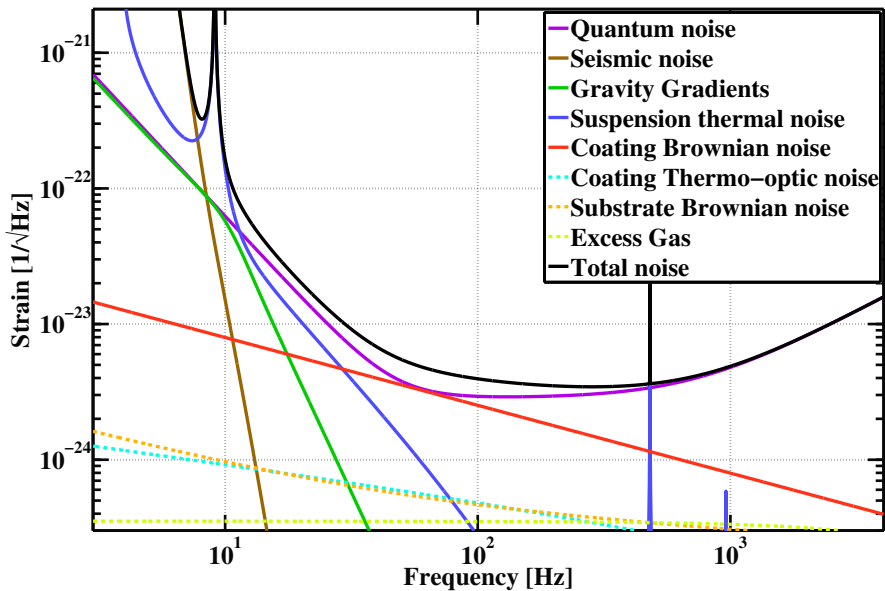


Classes of sources and searches

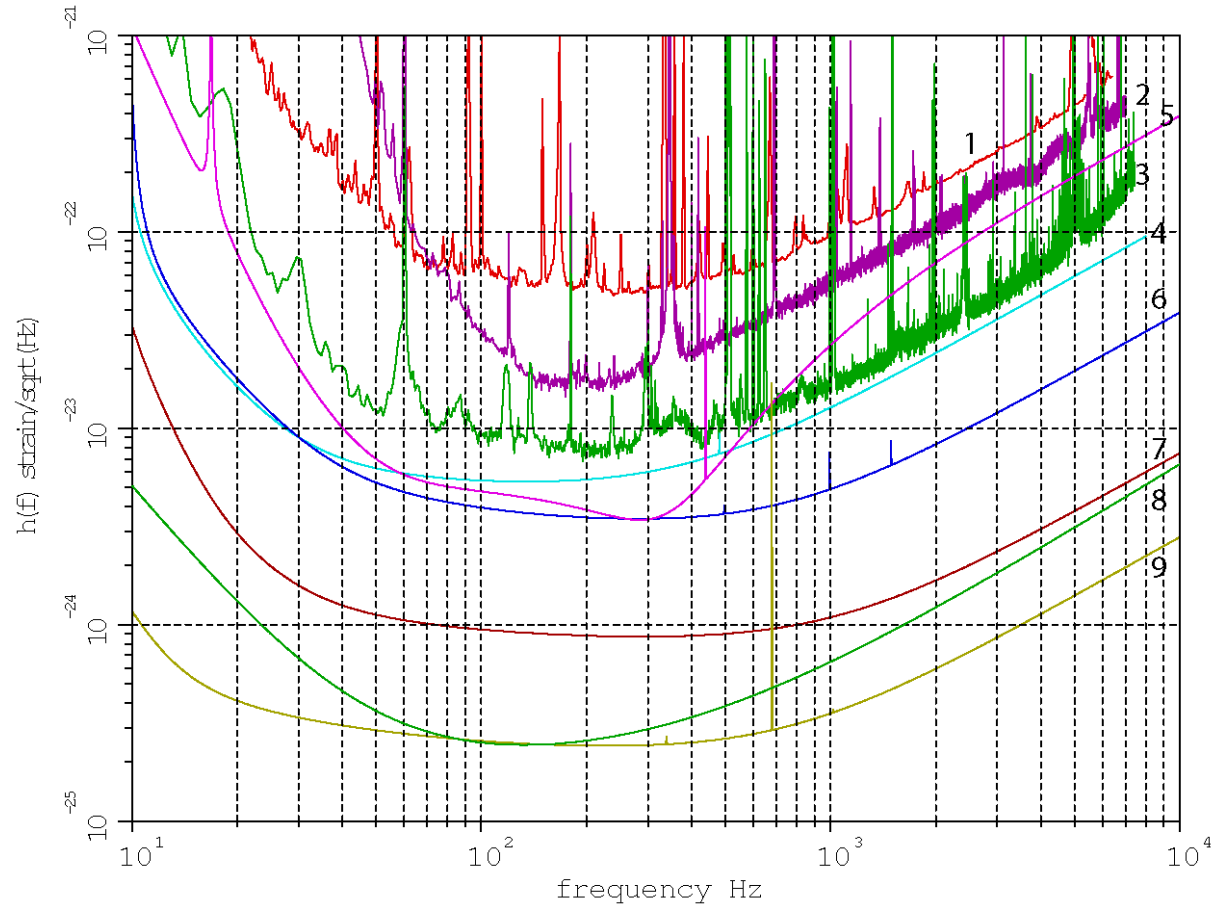
- **Compact binary inspiral: template search**
 - BH/BH
 - NS/NS and BH/NS
- **Low duty cycle transients: wavelets, T/f clusters**
 - Supernova
 - BH normal modes
 - Unknown types of sources
- **Triggered searches**
 - Gamma ray bursts
 - EM transients
- **Periodic CW sources**
 - Pulsars
 - Low mass x-ray binaries (quasi periodic)
- **Stochastic background**
 - Cosmological isotropic background
 - Foreground sources : gravitational wave radiometry



Advanced LIGO 190Mpc NS/NS noise budget



Evolution of gravitational strain sensitivity

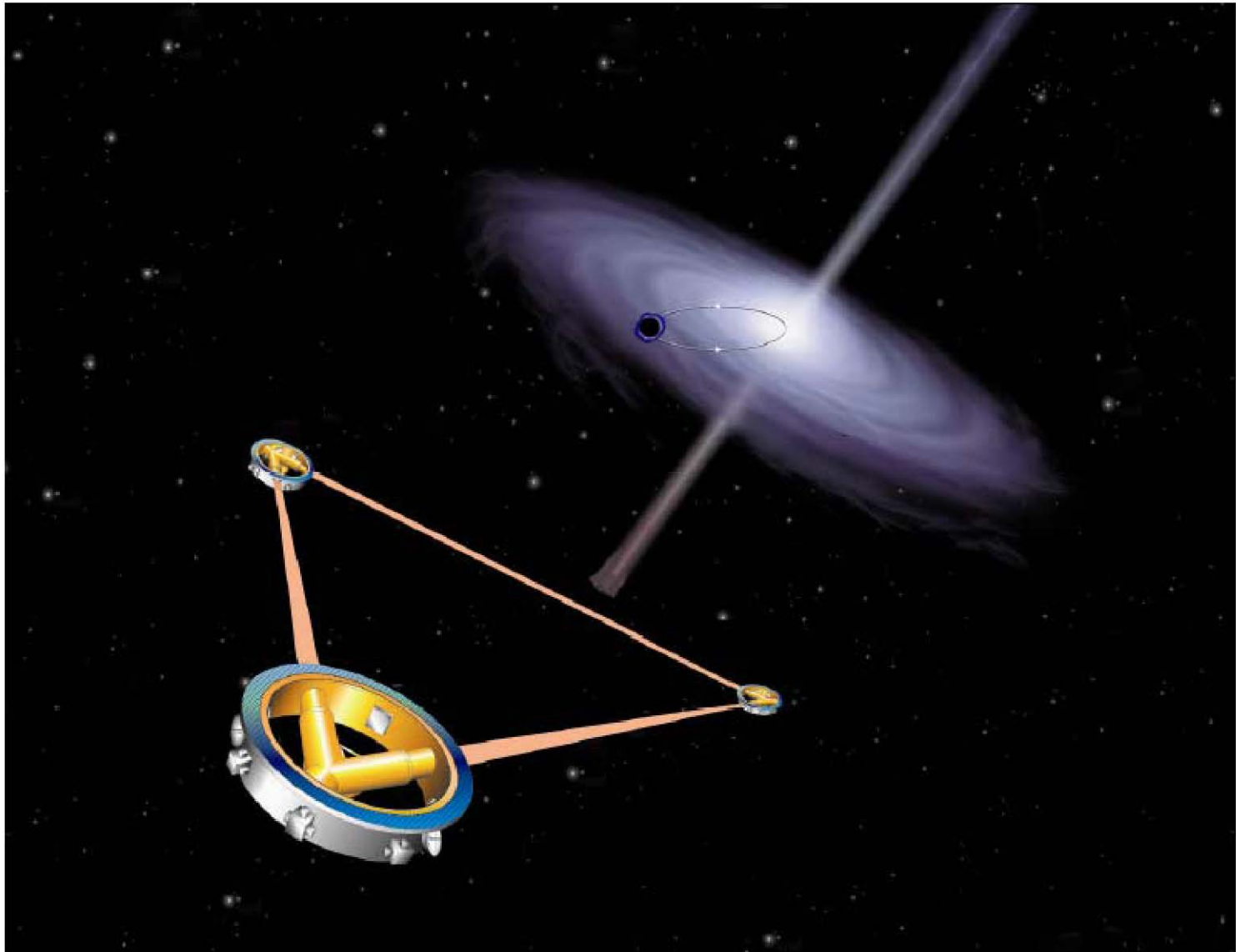


- 1 VIRGO 2009
- 2 Enhanced LIGO 2009
- 3 Advanced LIGO 65Mpc NS/NS 2015
- 4 Advanced LIGO 150Mpc NS/NS Low Power
- 5 Advanced VIRGO
- 6 Advanced LIGO 190Mpc NS/NS High Power
- 7 4km "Voyager" example 600Mpc NS/NS
- 8 Einstein telescope B
- 9 40km "Cosmic Explorer" example

Epoch	Estimated Run Duration	$E_{\text{GW}} = 10^{-2} M_{\odot} c^2$ Burst Range (Mpc)		BNS Range (Mpc)		Number of BNS Detections	% BNS Localized within	
		LIGO	Virgo	LIGO	Virgo		5 deg ²	20 deg ²
2015	3 months	40 – 60	–	40 – 80	–	0.0004 – 3	–	–
2016–17	6 months	60 – 75	20 – 40	80 – 120	20 – 60	0.006 – 20	2	5 – 12
2017–18	9 months	75 – 90	40 – 50	120 – 170	60 – 85	0.04 – 100	1 – 2	10 – 12
2019+	(per year)	105	40 – 80	200	65 – 130	0.2 – 200	3 – 8	8 – 28
2022+ (India)	(per year)	105	80	200	130	0.4 – 400	17	48



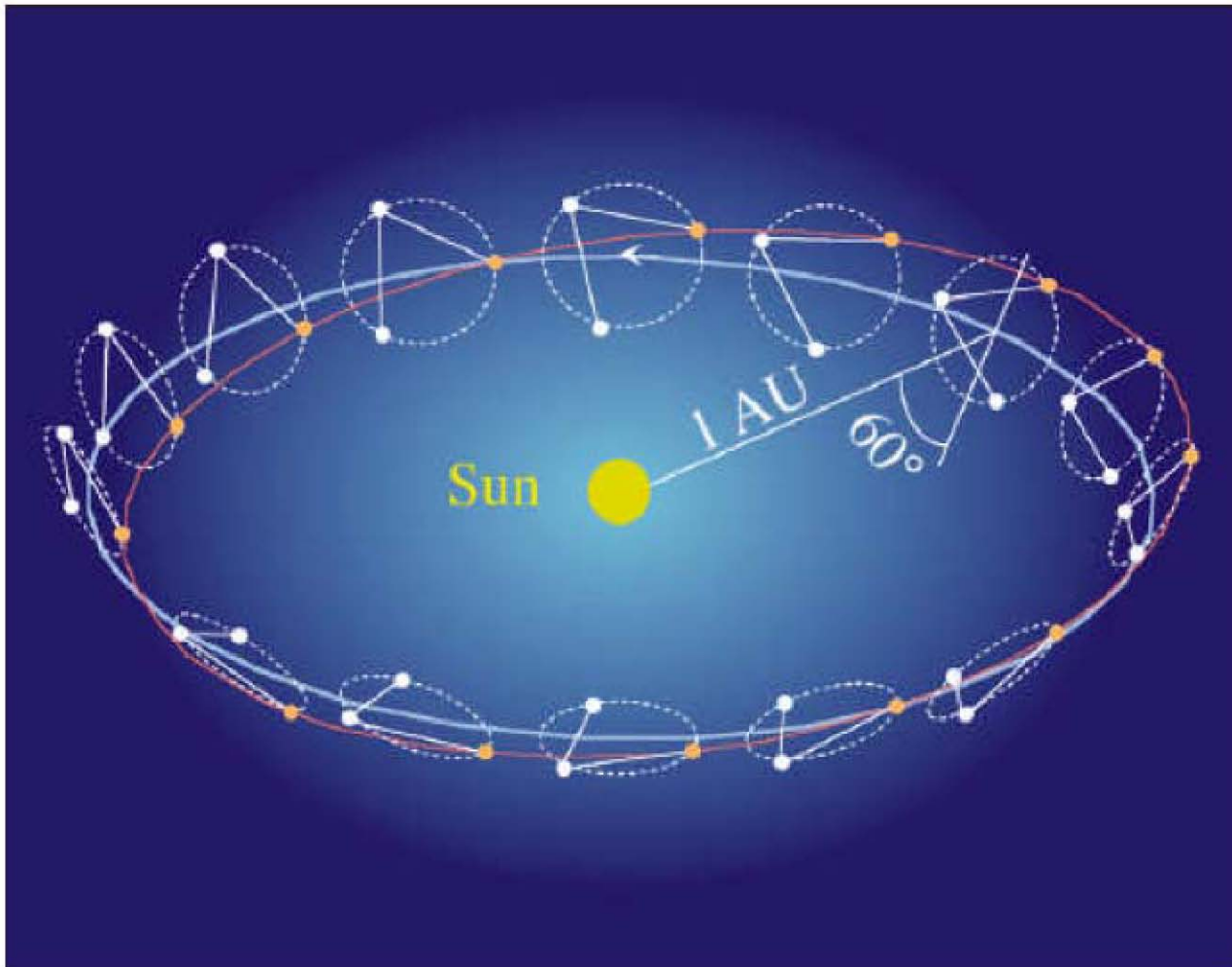
Mission Concept





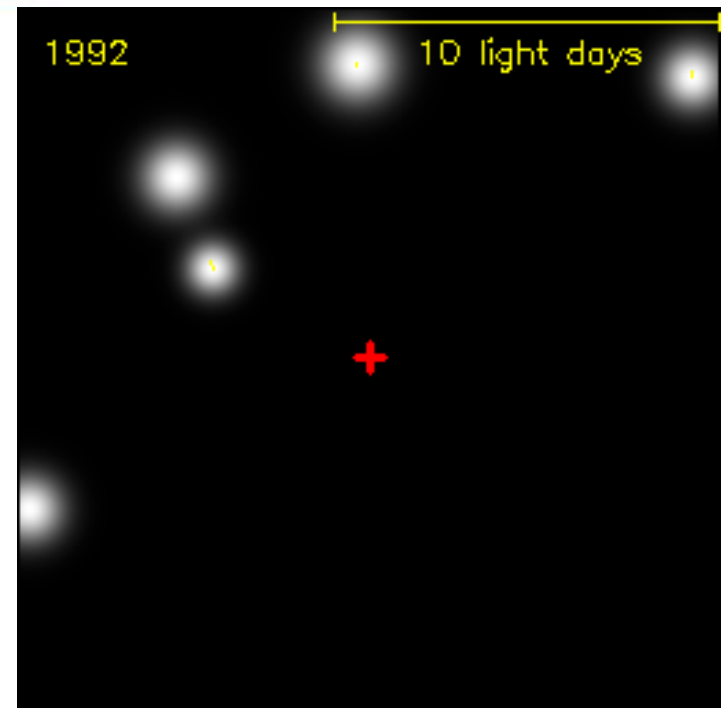
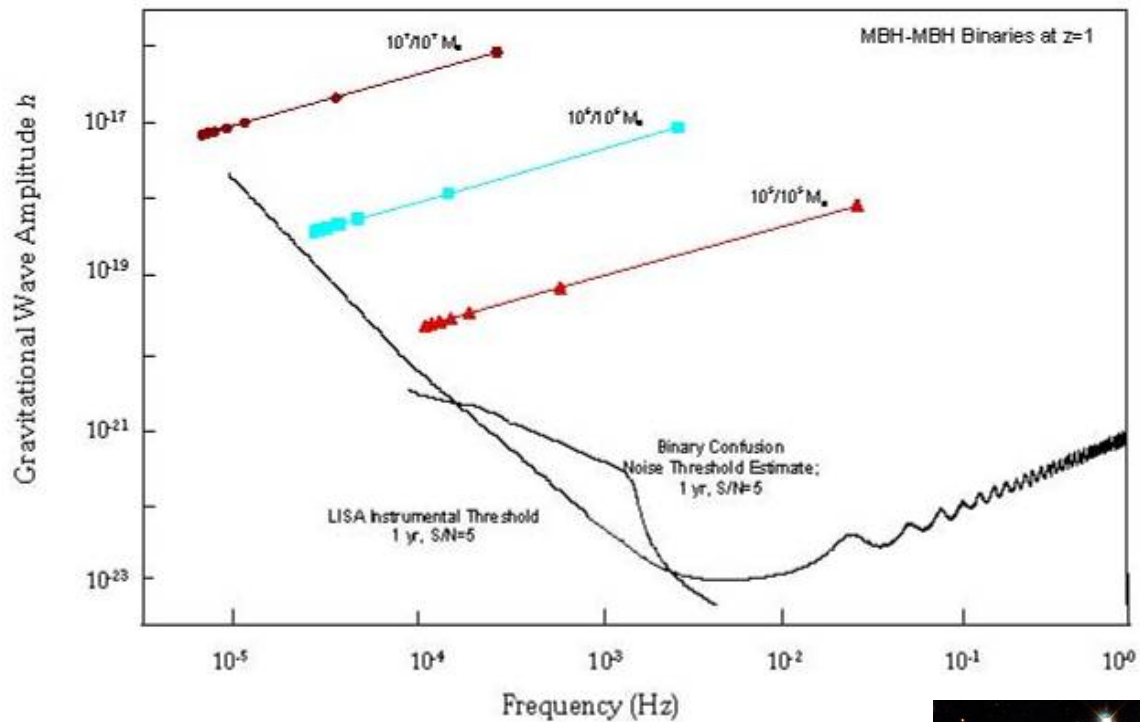
Spacecraft Orbits

- Spacecraft orbits evolve under gravitational forces only
- Spacecraft fly “drag-free” to shield proof masses from non-gravitational forces





Massive Black Holes in Merging Galaxies

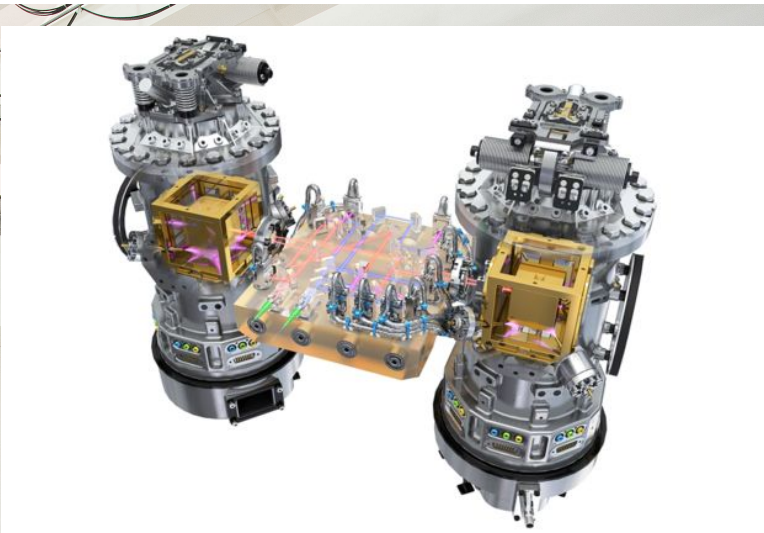
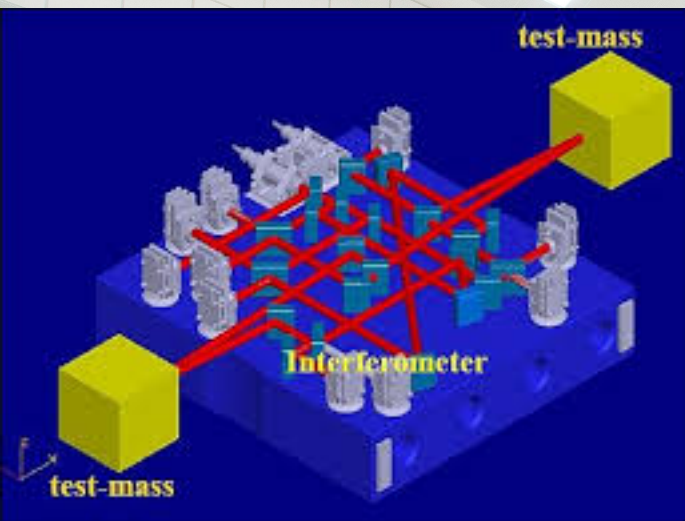


R. Genzel

Hubble Space Telescope



LISA Pathfinder



To be launched 12/02/2015

References

Traveling at the speed of thought

Daniel Kennefick

Princeton University Press 2007

Subtle is the Lord: The Science and Life of Albert Einstein

Abraham Pais

Clarendon Press 1982

Einstein in Berlin

Thomas Levenson

Random House 2003

